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Nuclear Structure of the Samarium Isotopes $^{152-154}$Sm Using models of IBM-2 and DDM

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Abstract: Samarium isotopes (Z=62) lie in the traditional rotational to transitional-spherical region that occurs at the range of deformed nuclei. Energy levels, Electromagnetic transitions B(E2), B(M1), mixing ratio $\delta(E2/M1)$ and monopole transitions $B(E0)$ for the selected transitions in $^{152\text{-}154}$Sm are calculated in the frame work of collective models proton-neutron Interacting boson model (IBM-2) and Dynamic deformation model (DDM). The results obtained for $^{152\text{-}154}$Sm are reasonably in good agreement with the known experimental results. The E0 quantities predicted by the two Models (IBM-2 and DDM) for 0_3^+ states were on the whole in poor agreement with experiment. Furthermore, the microscopic characters of the ground states and the excited 0^+ states in $Sm^{152\text{-}154}$ would appear to be sufficiently complex to eliminate any description of them in terms of the simple model. We noticed that most of the experimental and calculated values for the $X(E0/E2)$ ratio are small, which means that there is a small contribution of E0 transition on the life time of the 0^+ states.

Key words: Interacting boson model; dynamic deformation model; electromagnetic transition probabilities; mixing ratios; electric monopole transitions.

1. Introduction

Accurate information about the multipole moments of the low-lying excited states of ^{152}Sm and ^{154}Sm isotopes is of a particular interest, since these nuclei are in the transitional region between spherical and deformed nuclei. These spectra of states of the transitional nuclei fit neither a vibrational nor a rotational pattern, although they exhibit some characteristics of both. The low-lying excited states of ^{152}Sm and ^{154}Sm isotopes can be classified into ground, beta and gamma bands[1].

The phase transition from spherical to deformed shape which takes place in Sm isotopes has stimulated many authors to study theoretically and experimentally the area of these isotopes. The transition occurs between ^{150}Sm and ^{152}Sm. The former is a vibrational-like isotope while the latter is a rotational like one.

The isotopes Z=62 and N=90 to 92 are well abrupt changes in nuclear properties between almost spherical in N=82 to well deformed in N = 92 [2]. So Sm isotopes have provided a useful testing ground for nuclear structure calculations. Ref. [3] showed that the Sm isotopes lie in a transitional region from vibrational U(5) to the rotational limit SU(3) of the IBM. A description of such situation denoted X(5) as discussed in [4] , is a dynamical symmetry , albeit of unusual nature has also been developed , which describes at the critical point of a first order phase transition in two variables (vibrator to axial rotor transition region in atomic nuclei), is closely manifested empirically in ^{152}Sm [5].

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Ref. [6] used the interacting boson model (IBM) to interpret the ground state charge distributions and lowest 2^+ transition charge densities of the even $^{144-154}$Sm . Also in [7] used the effective IBM Hamiltonian to describe the low-lying energy spectra of several series of even-even isotopes. Applications are made to the Xe, Ba, Sm, Gd and U isotopes. In each case a single effective IBMA Hamiltonian containing at most six parameters reproduces some 50 experimental energies with an root mean square deviation of about 0.1MeV or less.

Ref. [8] studied the electric monopole transitions and structure of ^{150}Sm in IBM-2, and showed that the contour plot of the potential energy surfaces $V(\beta,\gamma)$ shows that ^{150}Sm is a spherical nucleus and has vibrational characters.

The aim of this work is to study the nuclear structure and electromagnetic transitions in $^{152-154}$Sm isotopes within the framework of IBM-2 and DDM and to compare the results with the available experimental data. In this work, predictions of the IBM-2 and DDM for the energy levels, electromagnetic transitions probabilities B(E2) and B(M1), multipole mixing ratios and monopole matrix elements in $^{152-154}$Sm isotopes are presented.

2.The Models

2.1 Interacting Boson Model

The Interacting Boson Model (IBM) [9-10] has been remarkably successful in describing the low-lying collective states in many medium to heavy even-even nuclei. The neutron-proton version of the interacting boson model (IBM-2), later suggested by [11] distinguish between neutron (ν) and proton (π) boson, is used in the present work, and a full description of the IBM-2 is found in ref. [12]. The Hamiltonian operator in IBM-2, which has been used to calculate the energy level and hence the gamma transitions matrix elements, has three parts, the first part is for proton bosons, the second is for neutron bosons whereas the last one is used to describe the interaction between different bosons:

$$H = H_\pi + H_\nu + V_{\pi\nu} \qquad (1)$$

The Hamiltonian generally used in the phenomenological calculations can be written as:

$$H = \varepsilon_d (n_{d\nu} + n_{d\pi}) + \kappa (Q_\nu . Q_\pi) + V_{\nu\nu} + V_{\pi\pi} + M_{\nu\pi} \qquad (2)$$

Where the dot denotes the scalar product. The first term represents the single-boson energy for neutron and proton, ε_d is the energy difference between *s*- and *d*-boson and $n_{d\rho}$ is the number of d-bosons, where ρ corresponds to π (proton) or ν (neutron) bosons. The second term denotes the main part of the boson-boson interaction, i.e., the quadrupole-quadrupole interaction between neutron and proton bosons with the strength κ. The quadrupole operator is:

$$Q_\rho = [d_\rho^{+} s_\rho + s_\rho^{+} d_\rho]^{(2)} + \chi_\rho [d_\rho^{+} d_\rho]^{(2)} \qquad (3)$$

Where χ_ρ determines the structure of the quadrupole operator and is determined empirically. The square brackets in eq.(3) denote the angular momentum coupling.

The terms $V_{\pi\pi}$ and $V_{\nu\nu}$, in eq. (2) correspond to the interaction between similar bosons, and there are sometimes included in order to improve its fits to experimental energy spectra. They are of the form:

$$V_{\rho\rho} = \frac{1}{2} \sum_{L=0,2,4} C_L^{\rho} ([d^{+}_{\rho} d^{+}_{\rho}]^{(L)} . [d_\rho d_\rho]^{(L)}). \qquad (4)$$

However, their effects are usually considered minor and often neglected [13].

The Majorana term, $M_{\pi\nu}$, shifts the states with mixed proton-neutron symmetry with respect to the totally symmetric ones. Since the little experimental information is known about such states with mixed symmetry, we did not attempt to fit the parameters appearing in eq. (2), but rather took constant values for the $^{152-154}$Sm isotopes (which contains three parameters ξ_1 , ξ_2 and ξ_3) may be written as:

$$M_{\pi\nu} = \frac{1}{2} \xi_2 ([s_\nu^{+} d_\pi^{+} - d_\nu^{+} s_\pi^{+}]^{(2)} . [s_\nu d_\pi - d_\nu s_\pi]^{(2)}) - \sum_{k=1,3} \xi_k ([d_\nu^{+} d_\pi^{+}]^{(k)} . [d_\nu d_\pi]^{(k)}). \qquad (5)$$

In t his work we give the Hamiltonian of IBM-2 and DDM terms of the formalism; and study the nuclear structures of $^{152-154}$Sm isotopes by these Hamiltonians.

2.2 Dynamic Deformation Model (DDM)

The dynamic deformation model has been developed over many years starting from the Paring Plus Quadrupole model (PPQ) of [14-15]. The DDM is an ambitious attempt to the collective spherical-transitional-deformed transitions and to span from the s-d shell to heavy nuclei using a microscopic theory of collective motion. No fitting parameters are required to obtain the data for a particular nucleus.

The detailed formalism and early results may be found in [16] and [17]. Here we give briefly the main aspects of the model. The theory can be divided into two main parts: a microscopic derivation of a collective Hamiltonian, and a numerical solution of the Hamiltonian. The microscopic Hamiltonian is composed of a demoralized Nilsson-type single particle plus pairing and has the form:

$$H = H_{av} + V_{res} \qquad (6)$$

Where

$$H_{av} = \frac{p^2}{2M} + \frac{1}{2}M\sum_{k=1}^{3}\omega_k^2\chi_k^2 + \hbar\omega_0\left[\upsilon_{ls}l.s + \upsilon_{ll}\left(l^2 - < l^2 >_N\right)\right] \qquad (7)$$

Combining all the various contributions together, the potential energy is written as:

$$V_{coll} = V_{DM} + \delta U + \delta V_{proj} + \delta E_{pair} \qquad (8)$$

Where δV_{proj} is a nine-dimensional projection correction introduced by [17]. The generalized cranking method is employed to derive the general expression for mass parameters $B_{\mu\nu}(\beta,\gamma)$ as used in the collective kinetic energy which can be written as:

$$T_{coll} = \frac{1}{2}\sum_{\mu\nu} B_{\mu\nu}\alpha_{\mu}^{\cdot}\alpha_{\mu}^{*} \qquad (9)$$

This kinetic energy function is quantized by Pauli method. The DDM code used for our calculation is a modified version of the latest DDM code which was developed for superheavy nuclei. The single particle levels and the configuration space (n = 0 to 8) employed in the present calculation, as well as the deformation definition, are identical to those of Kumar *et al*., [16].

3. Results and Discussion

3.1 Energy Levels

The isotopes chosen in this work are A=152, 154 due to the presence of the experimental data for the energy levels. We have $N_\pi = 6$, (12 protons outside the closed shell 50), and N_ν varies from 4 for ^{152}Sm to 5 for ^{154}Sm, measured from the closed shell at N=82 . While the parameters $\kappa, \chi_\rho, and\ \varepsilon_\rho$, as well as the Majorana parameters ξ_k, with k =1,2,3, were treated as free parameters and their values were estimated by fitting with the experimental values. In the calculation of $\delta(E2/M1)$, it is found that there is a great effect of the Majorano parameter ξ_2 on the value and sign of E2 and M1 matrix elements. The procedure was made by selecting the traditional value of the parameters and allowing one parameter to vary while keeping the others constant until the best fit with the experiment is obtained. This was carried out until one overall fit was obtained. The best values for the Hamiltonian parameters of IBM-2 are given in table 1. In the DDM there is no parameter fitted in the Hamiltonian except for Z (atomic number) and A (mass number).

Table 1: IBM-2 Hamiltonian parameters, all parameters in MeV units

Isotopes	ε	κ	χ_ν	χ_π	$\xi_1 = \xi_2$	ξ_3	$C_{L\nu}(L=0,2,4)$	$C_{L\pi}(L=0,2,4)$
^{152}Sm	0.43	-0.023	-0.8	-1.2	0.12	0.1	-0.5, 0.4, -0.8	-0.5, 0.4, -0.8
^{154}Sm	0.34	-0.039	-0.8	-1.2	0.12	0.1	-0.5, 0.4, -0.8	-0.5, 0.4, -0.8

A Concentration was made on the 2_1^+ to make a reasonable fits to the experimental data. A sample of experimental and theoretical values of energy levels is taken from table 2. It is noticed that a good agreement was obtained for the gamma and beta bands for ^{152}Sm. In addition, table 2 shows a comparison between the experimental and theoretical energy levels of the ground band in $^{152\text{-}154}$Sm isotopes. There is also an agreement between 2_1 and 4_1, but the DDM model is not able to predict the 6_1 and this may be due to the high spin of this state. Actually this has slim effects on the calculations of transitions probability.

The experimental energy ratio between first and second excited state is $R_{4/2} = E(4_1^+)/E(2_1^+)$ which has limiting values of 2 for quadrupole vibrator, 2.5 for gamma-soft and 3.33 for an ideal symmetric rotor. The ratio $R_{4/2}$ for ^{152}Sm equals 3.0004 for the experimental data, 3.048 for IBM-2 , and 3.304 for DDM results. The ratio $R_{4/2}$ for ^{154}Sm equal 3.256 for experimental data , 3.191 for IBM-2 and 3.154 for DDM. Therefore, from these ratio values, we deduced that the considered nuclei lie in the transitional region from gamma-soft a rotational shape O(6)- SU(3), and nearly close to the rotational limit.

Table 2: A comparison between the experimental and calculated energy levels for $^{152\text{-}154}$Sm in (MeV unit).

J_i^+	^{152}Sm			^{154}Sm		
	Exp. [18]	**IBM-2**	**DDM**	**Exp. [18]**	**IBM-2**	**DDM**
0(1)	0.0	0.0	0.0	0.0	0.0	0.0
2(1)	0.1218	0.123	0.114	0.082	0.0821	0.084
4(1)	0.366	0.375	0.387	0.267	0.262	0.265
0(2)	0.685	0.714	0.715	1.100	1.153	1.057
6(1)	0.706	0.711	0.585	0.549	0.602	0.631
2(2)	0.811	0.952	0.855	1.178	1.213	1.192
4(2)	1.023	1.001	1.039	1.371	1.421	1.430
2(3)	1.086	1.272	1.397	1.440	1.521	1.438
3(1)	1.234	1.003	1.559	1.539	1.557	1.541
4(3)	1.372	1.251	1.705	1.660	1.783	1.664
5(1)	-	1.432	1.888	-	1.980	2.321
6(2)	1.311	1.414	1.312	-	1.450	2.365
0(3)	1.080	1.139	1.473	1.202	1.201	1.450
2(4)	1.239	1.593	1.893	-	1.621	2.540
4(4)	-	2.211	2.210	1.662	1.783	1.982
8(1)	1.1250	1.290	1.254	0.903	1.021	0.973

4. Electromagnetic Transition Probability

4.1 Electric quadrupole Transition Probability

In IBM-2, the E2 , transition operator is given by:

$$T^{(E2)} = e_\pi Q_\pi + e_\nu Q_\nu \qquad (10)$$

Where Q_ρ is the same as in eq.(3) , e_π and e_ν are boson effective charges depending on the boson number N_ρ $(\rho = \pi$ or $\nu)$ and they can take any value to fit the experimental results $(B(E2;2_1^+ \to 0_1^+))$. The method is explained in ref. [19], and the effective charges calculated by this method for $^{152\text{-}154}$Sm isotopes were $e_\nu = 0.10\,eb$ and $e_\pi = 0.13eb$. Table 3 gives the electric transition probability B(E2).

The $B(E2;2_1^+ \to 0_1^+)$ and $B(E2;4_1^+ \to 2_1^+)$ values increased as neutron number increases toward the middle of the shell while the value of $B(E2;2_2^+ \to 2_1^+)$ has a small value because it contain mixtures of M1. The value of $B(E2;2_2^+ \to 0_1^+)$ is small because this transition is forbidden (from quasibeta band to ground state band) whereas the values of IBM-2 and DDM are in good agreement with the available experimental data.

In table 3, the electric quadrupole moment $Q(2_1^+)$ for the first excited state in $^{152\text{-}154}$Sm isotopes are well described indicating that the nucleus ^{152}Sm has a prolate shape in its first excited state.

Table 3: Electric Transition probability *B*(*E*2) for $^{152\text{-}154}$Sm isotopes in(e^2b^2 units)

$J_i^+ \to J_f^+$	^{152}Sm			^{154}Sm		
	Exp.(*)	**IBM-2**	**DDM**	**Exp.**	**IBM-2**	**DDM**
$2_1 \to 0_1$	0.670(15)	0.684	0.360	0.922(40)	0.913	0.940
$4_1 \to 2_1$	1.017(4)	0.980	0.932	1.186(39)	1.231	1.400
$6_1 \to 4_1$	1.179(33)	1.003	1.092	1.374(47)	1.393	1.350
$0_2 \to 2_1$	0.176(11)	0.120	0.454	0.235	0.250	0.235
$2_2 \to 0_1$	0.00456(34)	0.007	0.0252	0.060(14)	0.0071	0.013
$2_2 \to 2_1$	0.0258(26)	0.025	0.0752	0.012	0.014	1.970
$2_2 \to 4_1$	0.091(11)	0.077	0.165	0.024	0.028	0.054
$4_2 \to 2_1$	0.0035(35)	0.001	0.0049	-	0.027	0.012
$4_2 \to 4_1$	0.037(23)	0.041	0.0545	-	-	-
$2_3 \to 0_1$	0.0163(11)	0.027	0.662	0.013(3)	0.015	0.033
$2_3 \to 2_1$	0.0417(42)	0.051	0.1012	0.02	0.022	0.047
$2_3 \to 4_1$	0.0416(32)	0.003		-	0.0009	0.010
$4_3 \to 2_1$	0.0035(13)	0.008	0.0228	-	-	-
$4_3 \to 4_1$	0.037(13)	0.049	0.0211	-	-	-
$Q(2_1)$	-1.8(0.6)	-1.765	-1.64	-	-	-

* Experimental data are taken from references [18, 20,21,22,]

4.2 Magnetic Transition Probability B(M1)

After calculating the E2 matrix elements we look after the M1 matrix elements as fallows:

$$T^{(M1)} = \left[\frac{3}{4\pi}\right]^{1/2}(g_\pi L_\pi^{(1)} + g_\nu L_\nu^{(1)}) \qquad (11)$$

Where g_π, g_ν are the boson g-factors in units of μ_N and $L^{(1)} = \sqrt{10}(d^+ x \tilde{d})^{(1)}$. This operator can be written as:

The M1 operator is obtained by letting $l = 1$ in the single boson operator of the IBM-2 and can be written as:

$$T^{(M1)} = \left[\frac{3}{4\pi}\right]^{1/2}\left[\tfrac{1}{2}(g_\pi + g_\nu)(L_\pi^{(1)} + L_\nu^{(1)}) + \tfrac{1}{2}(g_\pi - g_\nu)(L_\pi^{(1)} - L_\nu^{(1)})\right] \qquad (12)$$

The first term on the right hand side of Eq.(12), is diagonal, and therefore, M1 transitions can be written as:

$$T^{(M1)} = 0.77\left[(d^+\tilde{d})_\pi^{(1)} - (d^+\tilde{d})_\nu^{(1)}\right](g_\pi - g_\nu) \qquad (13)$$

The direct measurement of B(M1) matrix elements is normally difficult, so the M1 strength of gamma transition may be expressed in terms of the multipole mixing ratio which can be written as [23]

$$\delta(E2/M1) = 0.835\,E_\gamma(MeV).\frac{< J_f|T(E2)|J_i >}{< J_f|T(M1)|J_i >} \quad \text{In eb/ } \mu_N \qquad (14)$$

Having fitted E2 matrix elements, one can use them to obtain M1 matrix elements and then the mixing ratio $\delta(E2/M1)$, compare them with the prediction of the model using the operator in (eq.9). The g_π *and* g_ν have to be estimated, if they were not measured in the case of Sm isotopes. The g factors may be estimated from the

$$g = g_\pi \frac{N_\pi}{N_\pi + N_\nu} + g_\nu \frac{N_\nu}{N_\pi + N_\nu}$$

Many relations could be obtained for a certain mass region and then the average g_π *and* g_π values for this region could be calculated. One of the experimental B(M1) and the relation above have been used to find out that $g_\pi - g_\nu = 0.53\mu N$. The estimated values of the parameter are $g_\pi = 0.84\,\mu_N$ *and* $g_\nu = 0.31\,\mu_N$.

experimental magnetic (μ) moment of the 2^+_1 state (μ=2g). In the phenomenological studies g_π and g_ν are treated as parameters, and kept constant for a whole isotope chain. The total g factor is defined by Sambataro *et. al*., [24] as:

(15)

These were used to calculate the mixing ratio $\delta(E2/M1)$. The ratios were calculated for some selected transitions and listed with the available experimental data in table 4. A good agreement is obtained between the calculated values by (IBM-2 and DDM) and the experimental data in sign and magnitude.

Table 4: Mixing Ratios for $^{152\text{-}154}$Sm

$J_i^+ \to J_f^+$	^{152}Sm			^{154}Sm		
	Exp. [18,24]	**IBM-2**	**DDM**	**Exp. [18,23]**	**IBM-2**	**DDM**
$2_2 \to 2_1$	8^{+9}_{-3}	9.50	10.8	56^{+130}_{-25}	34	0.054
$4_2 \to 4_1$	3(1)	5.102	4.0	-1.1	-0.970	0.320
$2_3 \to 2_1$	$-11^{+0.7}_{-0.8}$	-9.521	-24.3	0.8^{+15}_{-6}	0. 620	0.047
$3_1 \to 2_1$	-30^{+7}_{-12}	-8.0	-26.7	-7.5	-5.22	2.40
$3_1 \to 4_1$	$-12.2^{+1.2}_{-0.6}$	-11.32	-16.8	-	0.032	0.052
$2_3 \to 2_2$	-	6.6	19.6	-	0.007	0.003
$4_3 \to 4_1$	$-3^{+1}_{-2.4}$	-2.81	-9.6	-	0.0421	0.021
$3_1 \to 2_2$	-	7.41	6.9	-	0.0731	0.0982
$3_1 \to 4_2$	-	-7.51	-21.0	-	$1.2*10^{-4}$	$2.73*10^{-4}$
$4_3 \to 4_2$	-	4.462	9.70	-	$2.5*10^{-3}$	$3.72*10^{-3}$

The magnetic dipole transition probability is presented in table 5. It is worth mentioning here that there is no experimental data to compare with the theoretical results. The $B(M1; 2_3^+ \to 2_1^+)$ values decrease for two models , implying some collective effects.

The large B(M1) values in IBM-2 are due to the F-spin vector character of 2_3^+ state in $^{152\text{-}154}$Sm. The $B(M1;0_1^+ \to 1_1^+)$ is still sizable in $^{152\text{-}154}$Sm (increased with the increase of neutron number) because of the transition from ground state band to mixed symmetry state 1_1^+ in IBM-2 as well as in DDM.

Table 5: Magnetic Transition probability *B*(*M1*) for $^{152\text{-}154}$Sm isotopes in (μ^2_N units)

$J_i^+ \to J_f^+$	^{152}Sm		^{154}Sm	
	IBM-2	**DDM**	**IBM-2**	**DDM**
$2_2 \to 2_1$	0.07	$0.4*10^{-2}$	0.007	$4.081*^{1}0^{-4}$
$2_3 \to 2_1$	0.06	$0.282*^{1}0^{-2}$	0.020	$1.055*^{1}0^{-3}$
$0_1 \to 1_1$	1.460	1.367	1.810	$1.939*^{1}0^{-6}$
$3_1 \to 2_1$	0.262	$0.082*10^{-2}$	-	-
$3_1 \to 2_2$	0.082	$0.727*10^{-2}$	-	-
$3_1 \to 3_2$	0.087	$0.327*10^{-2}$	-	-

5. Monopole Transition Probability B(E0)

Monopole transitions (E0) are known to be pure penetration effect, where the transition is caused by an electromagnetic interaction between the nuclear charge and the atomic electron penetrating the nucleus. An E0 transition occurs between two states of the same spin and parity by transferring the energy and zero unit of angular momentum. Thus E0 has no competing gamma ray. These transitions are different from zero only in the case where the transition is accompanied by the nucleus surface change, for example in the nuclear models where the surface is assumed to be fixed E0 transitions are strictly forbidden. Electric monopole transitions can occur not only in $0^+ \to 0^+$ transition but also, in competition with gamma multipole transition and depending on transition selection rules may compete in any $\Delta I = 0$ decay such as $2^+ \to 2^+$. At transitions energies greater than $2m_o c^2$, monopole pair production is also possible.

The monopole transition operator T(E0) is given by [25]

$$T(E0) = \beta_{0\pi} d_\pi^+ . \tilde{d_\pi} + \beta_{0\nu} d_\nu^+ . \tilde{d_\nu} + \gamma_{0\nu} N_\pi + \gamma_{0\nu} N_\nu$$

$$B(E0; J_i - J_f) = e^2 R^4 \rho^2(E0) \quad J_i = J_f$$ E0 reduced transitions probability is written as in IBM-2 [25] (16)

Where e is the electronic effective charge, R is the nuclear radius and $\rho(E0)$ is the transition matrix element. However, there are only limited cases where $\rho(E0)$ can be measured directly. In most cases, we have to determine the intensity ratio of E0 to the competing E2 transition calling this as $X(E0/E2)$ value [26] which can be written as

$$X(E0/E2) = \frac{B(E0; J_i - J_f)}{B(E2; J_i - J_{f'})} \qquad (17)$$

Where $J_f = J_{f'}$ for $J_i \neq 0, and\ J_f = 0, J_{f'} = 2\ for\ J_i = 0.$

The $T^{(E0)}$ operator my be found by setting $l = 0$ on the IBM-2 operator [27]

$$\rho_{if}(E0) = \frac{Z}{R_0^2} \sum \tilde{\beta}_{0\rho} \left\langle f \left| d^+{}_\rho x d_\rho \right| i \right\rangle \qquad (18)$$

Where R_0=1.2A$^{1/3}$ fm and $\rho(E0)$ is a dimensionless quantity. The two parameters $\tilde{\beta}_{0\pi}, \tilde{\beta}_{0\nu}$ in equation (14) may be estimated by fitting in isotope shift, which is the difference in the square radius $\delta\langle r^2 \rangle$ between the neighboring isotopes in their ground state [28].i.e.

$$\delta\langle r^2 \rangle = \langle e | r^2 | e \rangle - \langle 0_1^+ | r^2 | 0_1^+ \rangle$$

$$\delta < r^2 >= \beta_{0n}\delta n_{d\pi} + \beta_{0\nu}\delta n_{d\nu} \text{.....(19)}$$

where $\delta n_{d\rho} = < e\left|d_\rho^+ d_\rho^{\sim}\right|e > - < 0\left|d_\rho^+ d_\rho^{\sim}\right|0 >$

In the case of Samarium isotopes, the measured isotopes shift [29] was used to find the parameters used in the IBM-2 calculations and they are $\tilde{\beta}_\pi = 0.045\ fm^2$ and $\tilde{\beta}_\nu = 0.031\ fm^2$ which produced the monopole matrix elements. The calculated and experimental values of $\rho(E0)$, $X(E0/E2)$, and the isomer shifts are given in tables 6,7and 8 respectively.

Table 6: Monopole matrix element ρ(E0) for $^{152\text{-}154}$Sm

$J_i^+ \rightarrow J_f^+$	^{152}Sm			^{154}Sm	
	Exp. [25]	**IBM-2**	**DDM**	**IBM-2**	**DDM**
$0_2 \rightarrow 0_1$	0.26(2)	0.311	0.365	0.432	0.563
$2_2 \rightarrow 2_1$	0.255(12)	0.402	0.363	0.532	0.672
$2_3 \rightarrow 2_1$	-	0.032	0.007	0.021	0.0117
$0_3 \rightarrow 0_1$	-	0.021	0.038	0.0371	0.0627
$0_3 \rightarrow 0_2$	0.261(36)	0.231	0.442	0.281	0.376

Table 7: X(E0/E2) for $^{152\text{-}154}$Sm

$J_i^+ \rightarrow J_f^+$	^{152}Sm			^{154}Sm	
	Exp. [25]	**IBM-2**	**DDM**	**IBM-2**	**DDM**
$0_2 \rightarrow 0_1$	0.7(0.1)	0.752	1.092	0.821	2.27
$2_2 \rightarrow 2_1$	4.5(0.5)	3.798	7.831	3.720	5.371
$4_2 \rightarrow 4_1$	6.6(2.10)	5.931	8.440	1.717	3.710

Table 8: isomer shift $\delta<r^2>$ for $^{152\text{-}154}$Sm

	^{152}Sm			^{154}Sm		
$\delta<r^2>$	**Exp. [30]**	**IBM-2**	**DDM**	**Exp. [30]**	**IBM-2**	**DDM**
	$18(4)*10^{-2}$	$16*10^{-2}$	$15.78*10^{-2}$	1.2(0.8)	2.51	12.9

The difference in the predications for X(5) and E(5) stemed rather from the ɣ-dependence behavior. To do so, [9], exploit the simplicity and flexibility of the Interacting Boson Approximation(IBA) Model .The model can be viewed as a huge truncation of shell model spaces to include only those configurations in which either two protons or two neutrons are coupled to spin or 2 (s and d bosons), respectively [30] .

6. Conclusions

In the present work, the energies of low-lying levels, E2 and M1 reduced transition probabilities for $^{152\text{-}154}$Sm isotopes were calculated in the framework of IBM-2 and DDM. The calculated energy levels of low lying states were well reproduced, though some discrepancies remain, especially in the high spin states. The transition probability B(E2) between the ground band, the quasi gamma band and quasi beta band states is also calculated. Good agreement with the experimental energy for the low lying levels was obtained. The mixing ratios were also calculated after calculating E2 and M1 matrix elements. All the experimental and theoretical mixing ratios for $^{152\text{-}154}$Sm isotopes indicate a small M1 component which means that in the band-mixing transitions, M1 components are almost forbidden However, an acceptable overall agreement between the experiment and theory was obtained as shown in table 2.

The E0 quantities predicted by the two Models (IBM-2 and DDM) for 0_3^+ states were on the whole in poor agreement with experiment. Furthermore, the microscopic characters of the ground states and the excited 0^+ states in $Sm^{152\text{-}154}$ would appear to be sufficiently complex to eliminate any description of them in terms of the simple model. We noticed that most of the experimental and calculated values for the $X(E0/E2)$ ratio are small, which means that there is a small contribution of E0 transition on the life time of the 0^+ states. There are two high values of $X(E0/E2)$ in transitions from 0_2^+ to 0_1^+ in $^{152\text{-}154}$Sm which means that this state decay mostly by the E0 is, the study of this state which gives information about the shape of the nucleus, because the E0 transitions matrix elements are connected strongly with the penetration of the atomic electron to the nucleus.

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