# NEW CRITERIA FOR OSCILLATION OF SECOND ORDER NONLINEAR DYNAMIC EQUATIONS WITH DAMPING ON TIME SCALES 

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#### Abstract

The oscillation of solutions of the second order nonlinear damped dynamic equation $\left(r(t) \psi(x(t)) x^{\Delta}(t)\right)^{\Delta}+$ $p(t) x^{\Delta}(t)+f(t, x(\tau(t)))=0$ on an arbitrary time scale $T$ is investigated. A generalized Riccati transformation is applied for the study of the Kamenev-type oscillation criteria for this nonlinear dynamic equation. Several new sufficient conditions for the oscillation of solutions are obtained to extend some known results in the literature.


KEYWORDS: Damped Delay Dynamic Equations, Oscillation Criteria, Time Scales

## INTRODUCTION

Much recent attention has been given to dynamic equations on time scales, we refer the reader to the landmark paper of S. Hilger [1]. Since then, several authors have expounded on various aspects of this new theory; see the survey paper by Agarwal, Bohner, O'Regan and Peterson [2]. A time scale $T$ is an arbitrary nonempty closed subset of the real numbers $R$. Thus, $R ; Z ; N ; N_{0}$, i.e., the real numbers, the integers, the natural numbers, and the nonnegative integers are examples of time scales. On any time scale $T$, we define the forward and backward jump operators by
$\sigma(t)=\inf \{s \in T, s>t\}, \rho(t)=\sup \{s \in T, s<t\}$.
A point $\mathrm{t} t \in T, t>\inf T$ is said to be left dense if $\rho(t)=t$, right dense if $t<\sup T$ and $\sigma(t)=t$, left scattered if $\rho(t)<t$, and right scattered if $\sigma(t)>t$. A function $f: T \rightarrow R$ is called rd-continuous provided that it is continuous at right dense points of $T$, and its left-sided limits exist (finite) at left-dense points of $T$. The set of rdcontinuous functions is denoted by $C_{r d}(T, R)$. By $C_{r d}^{1}(T, R)$, we mean the set of functions whose delta derivative belongs to $C_{r d}(T, R)$. In recent years, there has been much research activity concerning the oscillation and nonoscillation of solutions of various equations on time scales (see [5],[6],[7]). However, there are few results dealing the oscillation of solutions of delay dynamic equations on time scales [8-13]. Following this trend, we are concerned in this paper with oscillation
for the second-order nonlinear delay dynamic equations of the type

$$
\begin{equation*}
\left(r(t) \psi(x(t)) x^{\Delta}(t)\right)^{\Delta}+p(t) x^{\Delta}(t)+f(t, x(\tau(t)))=0 \tag{1.1}
\end{equation*}
$$

We assume that
$\left(\boldsymbol{H}_{\mathbf{1}}\right) r, \psi$ and $p$ are real-valued positive rd-continuous functions defined on $T$,
$0 \leq p(t) \leq 1$ and there are two positive constants $c_{1}, c_{2}$ such that $c_{1} \leq \psi(x(t)) \leq c_{2}$.
$\left(H_{2}\right) \tau: T \rightarrow R$ is strictly increasing, and $\tau(t) \leq t$ and $\tau \rightarrow \infty$ as $t \rightarrow \infty$
$\left(\boldsymbol{H}_{3}\right) f(t, u) \in C_{r d}(T \times R, R)$ satisfies $u f(t, u)>0$, for $u \neq 0$ and there exists a positive
rd-continuous function $q$ defined on $T$ such that $\left|\frac{f(t, u)}{u}\right| \geq q(t)$ for $u \neq 0$.

## 2. MAIN RESULTS

We need the following lemma for the proof of main results.
Lemma 2.1. Assume that $x(t)$ is a positive solution of Eq. (1.1) on $\left[t_{0}, \infty\right)_{T}$. If

$$
\begin{equation*}
\int_{0}^{\infty} \frac{e^{\frac{e}{-p(t)}\left(t, t_{0}\right)}}{r(t)} \Delta t=\infty \tag{2.1}
\end{equation*}
$$

and
$\int_{0}^{\infty} \tau(s) q(s) \Delta s=\infty$,
then there exists $t_{1} \in\left[t_{0}, \infty\right)_{T}$, such that
(i) $x^{\Delta}(t)>0,\left(r(t) \psi(x(t)) x^{\Delta}(t)\right)^{\Delta}<0$ for $t_{1} \in\left[t_{0}, \infty\right)_{T}$
(ii) $\frac{x(t)}{t}$ is decreasing.

Proof. Assume that $x(t)$ is a positive solution of Eq. (1.1) on $\left[t_{0}, \infty\right)_{T}$. Pick $t_{2} \in\left[t_{0}, \infty\right)_{T}$, such that $x(t)>0$ and $x(\tau(t))>0$ on $\left[t_{2}, \infty\right)_{T}$. Then without loss of generality we can take $x^{\Delta}(t)<0$ for all $t \geq t_{2} \geq t_{1}$. Now from (1.1) we have

$$
\begin{equation*}
\left(r(t) \psi(x(t)) x^{\Delta}(t)\right)^{\Delta}+p(t) x^{\Delta}(t)=-f\left(t, x^{\tau}(t)\right)<0 . \tag{2.3}
\end{equation*}
$$

Putting $y(t)=-r(t) \psi(x(t)) x^{\Delta}(t)$, then we can write (2.3) in the form

$$
-y^{\Delta}(t)-\frac{p(t)}{r(t) \psi(x(t))} \text { 回 } y(t)<0
$$

Thus
$y^{\Delta}(t)>-\frac{p(t)}{c_{1} r(t)}$ 回 $y(t)$
Therefore
$y(t)>y\left(t_{2}\right) e_{\frac{-p(t)}{c_{1} r(t)}}\left(t, t_{2}\right)$
i.e.
$-r(t) \psi(x(t)) x^{\Delta}(t)>-r\left(t_{2}\right) \psi\left(x\left(t_{2}\right)\right) x^{\Delta}\left(t_{2}\right) e_{\frac{-p(t)}{c_{1} r(t)}}\left(t, t_{2}\right)$

$$
x^{\Delta}(t)<\frac{r\left(t_{2}\right) \psi\left(x\left(t_{2}\right)\right) x^{\Delta}\left(t_{2}\right)}{\psi(x(t))} \frac{e_{\frac{-p(t)}{}\left(t, t_{2}\right)}^{c_{1} r(t)}}{r(t)}
$$

Then

$$
x(t)<x\left(t_{3}\right)+\frac{r\left(t_{2}\right) \psi\left(x\left(t_{2}\right)\right) x^{\Delta}\left(t_{2}\right)}{c_{1}} \int_{t_{3}}^{t} \frac{e_{-p(s)}^{c_{1} r(s)}}{c_{3}}\left(s, t_{2}\right) .
$$

By (2.1) we can see that $x(t) \rightarrow-\infty$ as $t \rightarrow \infty$ which contradicts the assumption $x(t)>0$ so $x^{\Delta}(t)>0$ and hence $\left(r(t) \psi(x(t)) x^{\Delta}(t)\right)^{\Delta}<0$ for all large $t$. Now to prove (ii), we define $V(t)=x(t)-t x^{\Delta}(t)$, if there is a $t_{3} \in$ $\left[t_{2}, \infty\right)_{T}$ such that $V(t)>0$. Suppose this is false, then $V(t)<0$ on $\left[t_{3}, \infty\right)_{T}$, thus

$$
\left(\frac{x(t)}{t}\right)^{\Delta}=\frac{t x^{\Delta}(t)-x(t)}{t \sigma(t)}=\frac{-V(t)}{t \sigma(t)}>0 .
$$

Hence $\frac{x(t)}{t}$ is strictly increasing on $\left[t_{3}, \infty\right)_{T}$. Pick $t_{4} \in\left[t_{3}, \infty\right)_{T}$, so that $\tau(t) \geq \tau\left(t_{4}\right)$ for all $t \geq t_{4}$. Then

$$
\frac{x(\tau(t))}{\tau(t)} \geq \frac{x\left(\tau\left(t_{4}\right)\right)}{\tau\left(t_{4}\right)}:=m>0
$$

$$
\begin{equation*}
x(\tau(t)) \geq m \tau(t) \tag{2.4}
\end{equation*}
$$

But from (1.1) we have

$$
\left(r(t) \psi(x(t)) x^{\Delta}(t)\right)^{\Delta}+p(t) x^{\Delta}(t)=-f\left(t, x^{\tau}(t)\right)<0
$$

i.e.

$$
\left(r(t) \psi(x(t)) x^{\Delta}(t)\right)^{\Delta} \leq-f\left(t, x^{\tau}(t)\right) \leq-q(t) x^{\tau}(t) \leq-\square m q(t) \tau(t)
$$

$$
r(t) \psi(x(t)) x^{\Delta}(t)-r\left(t_{4}\right) \psi\left(x\left(t_{4}\right)\right) x^{\Delta}\left(t_{4}\right) \leq-\square m \int_{t_{4}}^{t} q(s) \tau(s) \Delta s
$$

i.e.

$$
r\left(t_{4}\right) \psi\left(x\left(t_{4}\right)\right) x^{\Delta}\left(t_{4}\right) \geq m \int_{t_{4}}^{t} q(s) \tau(s) \Delta s
$$

This contradicts (2.2). Therefore $V(t)>0$ for all $t \in\left[t_{3}, \infty\right)_{T}$, and hence $x(t)>t x^{\Delta}(t)$. Also we see that

$$
\left(\frac{x(t)}{t}\right)^{\Delta}=\frac{t x^{\Delta}(t)-x(t)}{t \sigma(t)}=\frac{-V(t)}{t \sigma(t)}<0 .
$$

So $\frac{x(t)}{t}$ is strictly decreasing on $\left[t_{3}, \infty\right)_{T}$.
Theorem 2.2. Suppose that (2.1) holds. Furthermore, suppose that there exists a positive $\Delta$-differentiable function $\delta(t)$ such that, for all sufficiently large $t_{1} \in\left[t_{0}, \infty\right)_{T}$, one has

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \int_{t_{1}}^{t}\left\{\frac{\delta(s) q(s) \tau(s)}{\sigma(s)}-\frac{c_{2} s r(s) P^{2}(s)}{4 \delta(s) \sigma(s)}\right\} \Delta s=\infty \tag{2.5}
\end{equation*}
$$

where $P(t):=\delta^{\Delta}(t)-\frac{\delta(t) p(t)}{c_{2} r(t)}$. Then Eq. (1.1) is oscillatory.
Proof. Suppose to the contrary that $x(t)$ is a nonoscillatory solution of (1.1). We may assume that $x(t)>0, x(\tau(t))>$ $0 \forall t \in\left[t_{1}, \infty\right)_{T}$. The proof when $x(t)$ is eventually negative is similar. From Lemma 2.1 and Eq. (1.1) it follows that $x(\tau(t))>0, x^{\Delta}(t)>0,\left(r(t) \psi(x(t)) x^{\Delta}(t)\right)^{\Delta}<0 \quad \forall t \geq t_{1}$. Define
$\omega(t)=\delta(t) \frac{r(t) \psi(x(t)) x^{\Delta}(t)}{x(t)}$
Then
$\omega^{\Delta}(t)=\delta^{\Delta}\left[\frac{r \psi x^{\Delta}}{x}\right]^{\sigma}+\delta\left[\frac{r \psi x^{\Delta}}{x}\right]^{\Delta}$
$=\frac{\delta^{\Delta}}{\delta^{\sigma}} \omega^{\sigma}+\frac{\left(r \psi x^{\Delta}\right)^{\Delta}}{x^{\sigma}}-\delta \frac{r \psi\left(x^{\Delta}\right)^{2}}{x x^{\sigma}}$
From (1.1) and $\left(\boldsymbol{H}_{3}\right)$ we have
$\omega^{\Delta}(t) \leq \frac{\delta^{\Delta}}{\delta^{\sigma}} \omega^{\sigma}-\frac{\delta p}{r \psi} \frac{r \psi x^{\Delta}}{x^{\sigma}}-\delta q \frac{x^{\tau}}{x^{\sigma}}-\delta \frac{r \psi x^{\Delta}}{x^{\sigma}} \frac{r \psi x^{\Delta}}{r \psi x}$
Since $\sigma>t$, then in view of $\left(\boldsymbol{H}_{\mathbf{1}}\right)$, and $\left(r \psi x^{\Delta}\right)^{\Delta}<0$, we have
$\omega^{\Delta}(t) \leq \frac{\delta^{\Delta}}{\delta^{\sigma}} \omega^{\sigma}-\frac{\delta p}{c_{2} r} \frac{\left(r \psi x^{\Delta}\right)^{\sigma}}{x^{\sigma}}-\delta q \frac{x^{\tau}}{x^{\sigma}}-\frac{\delta}{c_{2} r} \frac{\left(r \psi x^{\Delta}\right)^{\sigma}}{x^{\sigma}} \frac{\left(r \psi x^{\Delta}\right)^{\sigma}}{x^{\sigma}} \frac{x^{\sigma}}{x}$
Since $\frac{x}{t}$ is strictly decreasing and $\tau<\sigma ; t<\sigma$, then $\frac{x^{\tau}}{x^{\sigma}} \geq \frac{\tau}{\sigma}, \quad \frac{x^{\sigma}}{x} \geq \frac{\sigma}{t}$, thus
$\omega^{\Delta}(t) \leq \frac{\delta^{\Delta}}{\delta^{\sigma}} \omega^{\sigma}-\frac{\delta p}{c_{2} r \delta^{\sigma}} \omega^{\sigma}-\frac{\delta q \tau}{\sigma}-\frac{\delta \sigma}{c_{2} r t\left(\delta^{\sigma}\right)^{2}}\left(\omega^{\sigma}\right)^{2}$
$=\frac{P}{\delta^{\sigma}} \omega^{\sigma}-\frac{\delta \sigma}{c_{2} r t\left(\delta^{\sigma}\right)^{2}}\left(\omega^{\sigma}\right)^{2}-\frac{\delta q \tau}{\sigma}$
i.e.
$\omega^{\Delta}(t) \leq\left[\frac{1}{\delta^{\sigma}} \sqrt{\frac{\delta \sigma}{c_{2} t r \delta^{\sigma}}} \omega^{\sigma}-\frac{P}{2} \sqrt{\frac{c_{2} t r}{\delta \sigma}}\right]^{2}+\frac{c_{2} \operatorname{tr} P^{2}}{4 \delta \sigma}-\frac{\delta q \tau}{\sigma}$
therefore
$\omega^{\Delta}(t) \leq \frac{c_{2} t r P^{2}}{4 \delta \sigma}-\frac{\delta q \tau}{\sigma}$
Integrating this inequality from $t_{2}$ to $t$, we get
$\int_{t_{2}}^{t}\left\{\frac{\delta(s) q(s) \tau(s)}{\sigma(s)}-\frac{c_{2} s r(s) P^{2}(s)}{4 \delta(s) \sigma(s)}\right\} \Delta s \leq \omega\left(t_{2}\right)-\omega(t) \leq \omega\left(t_{2}\right)$

This contradicts with (2.5), hence the proof is completed. By choosing $\delta(t)=1, t \geq t_{0}$ in Theorem 2.2 we have the following oscillation result.

Corollary 2.3. Assume that the assumptions of Theorem 2.2 hold and for all sufficiently large $t_{1}$,

$$
\lim _{t \rightarrow \infty} \sup \int_{t_{1}}^{t}\left\{\frac{q(s) \tau(s)}{\sigma(s)}-\frac{s p^{2}(s)}{4 c_{2} r(s) \sigma(s)}\right\} \Delta s=\infty
$$

Then every solution of Eq. (1.1) is oscillatory on $\left[t_{0}, \infty\right)_{T}$. By choosing $\delta(t)=t, \quad t \geq t_{0}$ in Theorem 2.2 we have the following oscillation result.

Corollary 2.4. Assume that the assumptions of Theorem 2.2 hold. Then every solution of Eq. (1.1) is oscillatory on $\left[t_{0}, \infty\right)_{T}$ provided that

$$
\begin{aligned}
& \lim _{t \rightarrow \infty} \sup \int_{t_{1}}^{t}\left\{\frac{s q(s) \tau(s)}{\sigma(s)}-\frac{c_{2} r(s) P^{2}(s)}{4 \sigma(s)}\right\} \Delta s=\infty \\
& \text { where } P(t)=1-\frac{t p(t)}{c_{2} r(t)}
\end{aligned}
$$

Now, we define the function space $\mathfrak{R}$ as follows: $H \in \Re$ provided $H$ is defined for $t_{0} \leq s \leq t, t, s \in$ $\left[t_{0}, \infty\right)_{T}, H(t, t)=0, H(t ; s) \geq 0$ and $H$ has a nonpositive continuous $\Delta$-partial derivative $H^{\Delta_{s}}(t, s) \geq 0$ with respect to the second variable and satisfies for some $h \in \mathfrak{R}$. The following theorem extends Theorem 2.2 of [5].

$$
H^{\Delta_{s}}(t, s)+H(t, s) \frac{P(s)}{\delta^{\sigma}(s)}=-\frac{h(t, s)}{\delta^{\sigma}(s)} \sqrt{H(t, s)}
$$

Theorem 2.5. Suppose that the assumptions of Theorem 2.2 hold. If there exists a positive functions $H, h \in \mathfrak{R}$ such that for all sufficiently large $t_{1} \in\left[t_{0}, \infty\right)_{T}$, one has then every solution of Eq. (1.1) is oscillatory on $\left[t_{0}, \infty\right)_{T}$.

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{H\left(t, t_{0}\right)} \int_{t_{0}}^{t}\left\{\frac{\delta(s) q(s) \tau(s)}{\sigma(s)} H(t, s)-\frac{c_{2} s r(s) h^{2}(t, s)}{4 \delta(s) \sigma(s)}\right\} \Delta s=\infty \tag{2.7}
\end{equation*}
$$

Proof. Suppose to the contrary that $x(t)$ is a nonoscillatory solution of (1.1). Then as in Theorem 2.2 we have $x(\tau(t))>$ $0, x^{\Delta}(t)>0,\left(r(t) \psi(x(t)) x^{\Delta}(t)\right)^{\Delta}<0 \forall t \geq t_{1} \in\left[t_{0}, \infty\right)_{T}$. We define $\omega(t)$ as in Theorem 2.2, then from (2.6) we have

$$
\begin{equation*}
\frac{\delta q \tau}{\sigma} \leq-\omega^{\Delta}+\frac{P}{\delta^{\sigma}} \omega^{\sigma}-\frac{\delta \sigma}{c_{2} r t\left(\delta^{\sigma}\right)^{2}}\left(\omega^{\sigma}\right)^{2} \tag{2.8}
\end{equation*}
$$

Multiplying (2.8) by $H(t, s)$ and integrating from $t_{1}$ to $t$, we get

$$
\begin{aligned}
& \int_{t_{1}}^{t} \frac{\delta(s) q(s) \tau(s)}{\sigma(s)} H(t, s) \Delta s \leq \int_{t_{1}}^{t} \omega^{\Delta}(s) H(t, s) \Delta s+\int_{t_{1}}^{t} \frac{P(s)}{\delta^{\sigma}(s)} \omega^{\sigma}(s) H(t, s) \Delta s \\
& -\int_{t_{1}}^{t} \frac{\delta(s) \sigma(s)}{c_{2} r(s) s\left(\delta^{\sigma}(s)\right)^{2}}\left(\omega^{\sigma}\right)^{2}(s) H(t, s) \Delta s \\
& =\omega\left(t_{1}\right) H\left(t, t_{1}\right)+\int_{t_{1}}^{t} H^{\Delta_{s}}(t, s) \omega^{\sigma}(s) \Delta s+\int_{t_{1}}^{t} \frac{P(s)}{\delta^{\sigma}(s)} \omega^{\sigma}(s) H(t, s) \Delta s
\end{aligned}
$$

$$
-\int_{t_{1}}^{t} \frac{\delta(s) \sigma(s)}{c_{2} r(s) s\left(\delta^{\sigma}(s)\right)^{2}}\left(\omega^{\sigma}\right)^{2}(s) H(t, s) \Delta s
$$

Thus

$$
\begin{aligned}
& \int_{t_{1}}^{t} \frac{\delta(s) q(s) \tau(s)}{\sigma(s)} H(t, s) \Delta s \leq \omega\left(t_{1}\right) H\left(t, t_{1}\right) \\
& -\int_{t_{1}}^{t}\left\{\frac{1}{\delta^{\sigma}(s)} \sqrt{\frac{\delta(s) \sigma(s) H(t, s)}{c_{2} r(s) s}} \omega^{\sigma}(s)+\frac{h(t, s)}{2} \sqrt{\frac{c_{2} r(s) s}{\delta(s) \sigma(s)}}\right\}^{2} \Delta s+\int_{t_{1}}^{t} \frac{c_{2} r(s) s h^{2}(t, s)}{4 \delta(s) \sigma(s)} \Delta s \\
& \leq \omega\left(t_{1}\right) H\left(t, t_{1}\right)+\int_{t_{1}}^{t} \frac{c_{2} r(s) s h^{2}(t, s)}{4 \delta(s) \sigma(s)} \Delta s
\end{aligned}
$$

i.e.

$$
\frac{1}{H\left(t, t_{1}\right)} \int_{t_{0}}^{t}\left\{\frac{\delta(s) q(s) \tau(s)}{\sigma(s)} H(t, s)-\frac{c_{2} s r(s) h^{2}(t, s)}{4 \delta(s) \sigma(s)}\right\} \Delta s \leq \omega\left(t_{1}\right)
$$

which contradicts (2.7) and hence the proof is completed.
Theorem 2.6. Suppose that (2.1) holds. Furthermore, suppose that there exists a function $g(t)$ such that $r(t) g(t)$ is a $\Delta-$ differentiable function and there exists a positive real rd-function $v(t)$ such that, for all sufficiently large $t_{1} \in$ $\left[t_{0}, \infty\right)_{T}$, one has

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \int_{t_{1}}^{t}\left\{E(s)-\frac{B^{2}(s)}{4 A(s)}\right\} \Delta s=\infty \tag{2.9}
\end{equation*}
$$

where $A(t):=\frac{v^{\sigma}(t) t}{c_{2} r(t) \sigma(t) v^{2}(t)}, \quad B(t):=\frac{v^{\Delta}(t)}{v(t)}-\frac{t v^{\sigma}(t) p(t)}{c_{2} r(t) \sigma(t) v(t)}+\frac{2 t v^{\sigma}(t) g(t)}{c_{2} \sigma(t) v(t)}$, and $\quad E(t):=\frac{v^{\sigma}(t) q(t) \tau(t)}{\sigma(t)}+$ $\frac{\operatorname{tr}(t) v^{\sigma}(t) g^{2}(t)}{c_{2} \sigma(t)}-\frac{t p(t) v^{\sigma}(t)}{c_{2} \sigma(t)}-v^{\sigma}(t)[r(t) g(t)]^{\Delta}$. Then Eq. (1.1) is oscillatory.

Proof. Suppose to the contrary that $x(t)$ is a nonoscillatory solution of (1.1). We may assume that $x(t)>0, x(\tau(t))>$ $0 \forall t \in\left[t_{1}, \infty\right)_{T}$. The proof when $x(t)$ is eventually negative is similar. From Lemma 2.1 and Eq. (1.1) it follows that $x(\tau(t))>0, x^{\Delta}(t)>0,\left(r(t) \psi(x(t)) x^{\Delta}(t)\right)^{\Delta}<0 \quad \forall t \geq t_{1} \in\left[t_{0}, \infty\right)_{T}$. Define

$$
u(t)=v(t)\left[\frac{r(t) \psi(x(t)) x^{\Delta}(t)}{x(t)}+r(t) g(t)\right]
$$

Then

$$
\begin{aligned}
& u^{\Delta}(t)=v^{\Delta}\left[\frac{r \psi x^{\Delta}}{x}+r g\right]+v^{\sigma}\left[\frac{r \psi x^{\Delta}}{x}+r g\right]^{\Delta} \\
& =\frac{v^{\Delta}}{v} u+v^{\sigma} \frac{\left(r \psi x^{\Delta}\right)^{\Delta}}{x^{\sigma}}-v^{\sigma} \frac{r \psi\left(x^{\Delta}\right)^{2}}{x x^{\sigma}}+v^{\sigma}[r g]^{\Delta}
\end{aligned}
$$

From (1.1) and ( $\boldsymbol{H}_{3}$ ) we have
$u^{\Delta}(t) \leq \frac{v^{\Delta}}{v} u-\frac{v^{\sigma} p}{r \psi} \frac{r \psi x^{\Delta}}{x} \frac{x}{x^{\sigma}}-v^{\sigma} q \frac{x^{\tau}}{x^{\sigma}}-\frac{v^{\sigma}}{r \psi} \frac{x}{x^{\sigma}}\left(\frac{u}{v}-r g\right)^{2}+v^{\sigma}[r g]^{\Delta}$
Using Lemma 2.1 and ( $\boldsymbol{H}_{\mathbf{1}}$ ), we get
$u^{\Delta}(t) \leq \frac{v^{\Delta}}{v} u-\frac{v^{\sigma} p t}{c_{2} r \sigma}\left(\frac{u}{v}-r g\right)-\frac{v^{\sigma} q \tau}{\sigma}-\frac{v^{\sigma} t u^{2}}{c_{2} r \sigma v^{2}}+\frac{2 v^{\sigma} t u}{c_{2} \sigma v}-\frac{v^{\sigma} t r g^{2}}{c_{2} \sigma}+v^{\sigma}[r g]^{\Delta}$
i.e.
$u^{\Delta}(t) \leq \frac{v^{\Delta}}{v} u-\frac{v^{\sigma} p t}{c_{2} r \sigma v} u+\frac{v^{\sigma} p t g}{c_{2} \sigma}-\frac{v^{\sigma} q \tau}{\sigma}-\frac{v^{\sigma} t}{c_{2} r \sigma v^{2}} u^{2}+\frac{2 v^{\sigma} t u}{c_{2} \sigma v}-\frac{v^{\sigma} t r g^{2}}{c_{2} \sigma}+v^{\sigma}[r g]^{\Delta}$
$=-A u^{2}+B u-E=-\left[\sqrt{A} u-\frac{B}{2 \sqrt{A}}\right]^{2}+\frac{B^{2}}{4 A}-E$
i.e.
$u^{\Delta}(t) \leq \frac{B^{2}(t)}{4 A(t)}-E(t)$
By integrating the above inequality, we obtain
$\int_{t_{2}}^{t}\left\{E(s)-\frac{B^{2}(s)}{4 A(s)}\right\} \Delta s \leq u\left(t_{2}\right)-u(t) \leq u\left(t_{2}\right)$
This contradicts (2.9), and the proof is completed. We can get the following result by choosing $v(t)=1$ in Theorem 2.6.

Corollary 2.7. Assume that the assumptions of Theorem 2.5 hold and for some $t_{1}$ sufficiently large. We have

$$
\begin{aligned}
& \lim _{t \rightarrow \infty} \sup \int_{t_{1}}^{t}\left\{E(s)-\frac{B^{2}(s)}{4 A(s)}\right\} \Delta s=\infty \\
& \text { where } A(t):=\frac{t}{c_{2} r(t) \sigma(t)}, B(t):=-\frac{\operatorname{tp(t)}}{c_{2} r(t) \sigma(t)}+\frac{2 t g(t)}{c_{2} \sigma(t)^{\prime}} \text { andE }(t):=\frac{q(t) \tau(t)}{\sigma(t)}+\frac{\operatorname{tr}(t) g^{2}(t)}{c_{2} \sigma(t)}-\frac{\operatorname{tp(t)g(t)}}{c_{2} \sigma(t)}-[r(t) g(t)]^{\Delta} .
\end{aligned}
$$

Then Eq. (1.1) is oscillatory.

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