# Pressure dependence of elastic constants for intermetallic compounds 

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(Received 23 Nov, 2013, Accepted 13 Feb, 2014)


#### Abstract

Elastic properties of solids are closely related to many fundamental properties such as EOS, thermal expansion, interatomic potentials etc. The relatively strong covalent bonding between the constituents of intermetallic compounds like $\mathrm{LiBC}, \mathbf{M g B}_{2}$ and $\mathrm{TiB}_{2}$ gives high Bulk Modulus in the range $\mathbf{1 5 0 - 2 2 0} \mathbf{G P a}$ at high melting point nearly $\mathbf{3 0 0 0 K}$. This behaviour makes these compounds very interesting for studies under high pressure and high temperature. In the present study, pressure-volume relationship at different temperatures for three intermetallic compounds is given. For the low compressions the increase in pressure is low but the rate of increase of pressure increases faster as we approach the high compression region. At higher temperature, the bulk modulus decreases and, therefore, the pressure required for producing same amount of compression becomes less and less.


Keywords: Intermetallic compounds, pressure, Bulk Modulus, Isotherm, parameter.

## INTRODUCTION

The variations of thermoelastic properties of solids at high temperature can be understood by using adequate form of EOS and by studying the thermal emissivity, bulk modulus and other properties ${ }^{1-3}$. The pressure-volume relationship at constant temperature of the mineral gives isothermal compressibility. The Murnaghan's equation as well as the Birch Murnaghan equation needs to be modified if we have to predict the pressure-volume relationship at elevated temperatures. Such modifications should be introduced in order to incorporate the effect of thermal pressure.
In order to understand the thermoelastic behaviour of solids adequately under the effect of pressure and at high temperatures, we have to develop models for estimating the variations of Bulk Modulus ${ }^{4}$. The pressure dependence of elastic constants has attracted the attention of theoretical as well as experimental workers because of their requirement in geophysical and geochemical problems. It is desirable to investigate high derivative thermoelastic properties and Equation of State (EOS) at high pressure and high temperature. These properties can be determined with the help of an EOS having a fundamental basis.

## MATERIAL AND METHODS

The general form of the strain in terms of volume ratio $\frac{V}{V_{0}}$ where $V_{0}$ is volume at $\mathrm{P}=0$ is defined as ${ }^{5}$ :

$$
\begin{equation*}
f=\frac{1}{n}\left(x^{-n / 3}-1\right) \tag{1}
\end{equation*}
$$

Where $x=\frac{v}{V_{0}}$ and $n=2$ for Eulerian strain
The reference state for Eulerian strain is deformed state.
Within the framework of strain theory developed by Stracey ${ }^{6}$ free energy is expanded as:

$$
\begin{equation*}
H=H_{0}+C_{2} f^{2}+C_{3} f^{3} \tag{2}
\end{equation*}
$$

Where; $H_{0}=$ free energy for reference state; $C_{2}$ and $C_{3}$ are coefficients in power series expansion; $f$ is parameter as defined by equation 1.
The pressure-volume relationship along isotherms at different temperatures is obtained using Birch Murnaghan third order EOS ${ }^{7 \& 8}$ which is derived from Eulerian strain theory and given as:

$$
\begin{equation*}
p=\frac{3}{2} K_{0}\left(x^{-7 / 3}-x^{-5 / 3}\right)\left[1+\frac{3}{4}\left(K_{0}^{\prime}-4\right)\left(x^{-2 / 3}-1\right)\right] \tag{3}
\end{equation*}
$$

Where $K_{0}, K_{0}^{\prime}$ are the values of isothermal Bulk Modulus and its first pressure derivative $\frac{d K}{d P}$ at $P=0$.
The values of Bulk Modulus $\mathrm{K}_{0}(\mathrm{~T})$ at different temperatures are obtained as:

$$
\begin{equation*}
K_{0}(T)=K_{0}(0)+A_{1} T-A_{2} T^{2}+A_{3} T^{3} \tag{4}
\end{equation*}
$$

Where $A_{1}=3.94 \times 10^{-3}, A_{2}=6.5241 \times 10^{-5}$ and $A_{3}=2.7048 \times 10^{-8}$ and
$K_{0}(0)=170.3 \mathrm{GPa}$ for $\mathrm{LiBC}, 151 \mathrm{GPa}$ for $M g B_{2}$ and 213 GPa for $\mathrm{Ti} B_{2}$
The change in $K_{0}^{\prime}$ with the temperature for different materials has been studied by Upadhyaya and Sharma ${ }^{10}$. It has been found that the rate of change of $K_{0}^{\prime}$ with temperature depends on the value of Bulk Modulus. Depending upon their method we have estimated:

$$
\begin{aligned}
\frac{d K_{0}^{\prime}}{d T} & =0.3 \times 10^{-3} \mathrm{~K}^{-1} \text { for } \mathrm{LiBC} \text { and } \mathrm{Mg} B_{2} \\
\& & \frac{d K_{0}^{\prime}}{d T}
\end{aligned}
$$

The input parameter $K_{0}$ and $K_{0}^{\prime}$ are thus obtained for $\mathrm{LiBC}, \mathrm{MgB}_{2}$ and $\mathrm{TiB}_{2}$ are reported in Table 1.
Bina and Helffrich ${ }^{9}$ have formulated a third order Eulerian strain equation of state for elastic constants as:

Where

$$
\begin{gather*}
C_{i j}=C_{i j}^{0}(1+f)^{5 / 2}\left[1-f\left\{5-3\left(\frac{d C_{i j}}{d P}\right)_{0}\left(\frac{K_{0}}{C_{i j}}\right)\right\}\right]  \tag{5}\\
f=\frac{1}{2}\left[\left(\frac{V_{0}}{V}\right)^{2 / 3}-1\right] \tag{6}
\end{gather*}
$$

It has been assumed that the relative changes in elastic constants and bulk modulus are i.e.

$$
\begin{equation*}
\frac{1}{C_{i j}^{0}}\left(\frac{d C_{i j}}{d P}\right)_{0}=\frac{1}{K_{0}}\left(\frac{d K}{d P}\right)_{0} \tag{7}
\end{equation*}
$$

Using eq. (5) and (7), we get:

$$
\begin{equation*}
\frac{C_{i j}}{C_{i j}^{0}}=(1+f)^{5 / 2}\left\{1-f\left[5-3\left(\frac{d K}{d P}\right)_{0}\right]\right\} \tag{8}
\end{equation*}
$$

Here

$$
C_{i j} \text { is } C_{i j}(T, P) \text { and } C_{i j}^{0}=C_{i j}(T, 0)
$$

Equation ${ }^{8}$ can be used to predict pressure dependence of elastic constants at a given temperature for materials provided values of $K_{0}^{\prime}$ is known at that temperature
Upadhyay and Sharma ${ }^{10}$ calculated pressure-volume relationship at different temperatures using eq. (3) for three intermetallic compounds viz. $\mathrm{LiBc}, \mathrm{MgB}_{2}$ and $\mathrm{TiB}_{2}$. The input data of $K_{0}, K_{0}^{\prime}$ used in the calculations are given in Table 1.
At $\mathrm{P}=0$ for $\mathrm{V}=\mathrm{V}_{0}$, we can determine the value of elastic constants at different temperatures using the relation ${ }^{11}$.

$$
\begin{equation*}
C_{i j}(T)=C_{i j}(0)+A_{1} T-A_{2} T^{2}+A_{3} T^{3} \tag{9}
\end{equation*}
$$

## RESULTS AND DISCUSSION

The pressure-volume results of $\mathrm{LiBc}, \mathrm{MgB}_{2}$ and $\mathrm{TiB}_{2}$ are reported in Table 2, Table 3 and Table 4 respectively.
Values of volume ratio represented by $\frac{V}{V_{0}}$ are infact $\frac{V(T, P)}{V(T, 0)}$ i.e. Isothermal compression.
At each temperature T, we have $\frac{V(T, P)}{V(T, 0)}=1$ at $P=0$
As the volume is decreased, pressure increases.
In the starting i.e. for low compressions ( $\frac{V}{V_{0}}$ from 1 to 0.9 ) the increase in pressure is slow but the rate of increase of pressure increases and becomes faster and faster as we approach the high compression region. This explains why the pressure-volume relationships are nonlinear. At higher temperature, the bulk modulus decreases and therefore the pressure required for producing the same amount of compression at higher temperature becomes less and less. Values of elastic constants thus determined at elevated temperatures $\mathrm{T}=0 \mathrm{~K}, 300 \mathrm{~K}, 600 \mathrm{~K}$ and 900 K have been used in equation (8) to find elastic constants at elevated pressures along different isothermals.
The results of $C_{i j}(T, P)$ for $\mathrm{LiBC}, \mathrm{MgB}_{2}$ and $\mathrm{TiB}_{2}$ are reported in table 2, table 3 and table 4 respectively.

Table 1: Input values of $K_{0}(\mathbf{G P a})$ and $K_{0}^{\prime}$ for $\mathrm{LiBC}, \mathrm{MgB}_{2} \& \mathrm{TiB}_{2}$

| $\mathbf{T}(\mathbf{K})$ | $\mathbf{L i B C}_{2}$ |  | $\mathbf{M g B}_{\mathbf{2}}$ |  | $\mathbf{T i B}_{\mathbf{2}}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\left.\boldsymbol{K}_{\mathbf{0}} \mathbf{( G P a}\right)$ | $\boldsymbol{K}_{\mathbf{0}}^{\prime}$ | $\boldsymbol{K}_{\mathbf{0}} \mathbf{( G P a )}$ | $\boldsymbol{K}_{\mathbf{0}}^{\prime}$ | $\boldsymbol{K}_{\mathbf{0}}(\mathbf{G P a})$ | $\boldsymbol{K}_{\mathbf{0}}^{\prime}$ |
| 0 | 170.3 | 3.76 | 151 | 4.0 | 213 | 2.10 |
| 300 | 166.3 | 3.85 | 147 | 4.09 | 200 | 2.16 |
| 600 | 155 | 3.94 | 135.7 | 4.18 | 197.7 | 2.22 |
| 900 | 140.3 | 4.03 | 121.4 | 4.27 | 183.4 | 2.28 |

Table 2: Values of elastic constants at different pressures and temperatures for LiBC.

|  | 0 K |  |  |  | 300 K |  |  |  | 600 K |  |  |  | 900 K |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $V / V_{0}$ | $\begin{gathered} \hline \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ | $\begin{gathered} \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ | $\begin{gathered} \hline \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ | $\begin{gathered} \hline \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ |
| 1 | 0 | 638 | 117 | 123 | 0 | 634 | 113 | 119 | 0 | 622 | 102 | 108 | 0 | 608 | 87.4 | 93.4 |
| 0.95 | 9.60 | 770 | 141 | 148 | 9.42 | 768 | 137 | 144 | 8.80 | 757 | 124 | 131 | 7.98 | 743 | 107 | 114 |
| 0.90 | 21.8 | 933 | 171 | 180 | 21.5 | 935 | 167 | 175 | 20.1 | 925 | 152 | 161 | 18.3 | 911 | 131 | 140 |
| 0.85 | 37.5 | 1136 | 208 | 219 | 36.9 | 1142 | 203 | 214 | 34.7 | 1132 | 186 | 197 | 31.7 | 1119 | 161 | 172 |
| 0.80 | 57.6 | 1391 | 255 | 268 | 56.9 | 1401 | 250 | 263 | 53.7 | 1394 | 229 | 242 | 49.2 | 1382 | 199 | 213 |
| 0.75 | 83.7 | 1713 | 314 | 330 | 83.2 | 1732 | 309 | 325 | 78.6 | 1727 | 283 | 300 | 72.2 | 1716 | 247 | 264 |
| 0.70 | 118 | 2129 | 390 | 410 | 118 | 2157 | 384 | 405 | 112 | 2157 | 354 | 374 | 103 | 2148 | 309 | 330 |
| 0.65 | 164 | 2674 | 490 | 515 | 164 | 2715 | 484 | 510 | 156 | 2721 | 446 | 472 | 145 | 2716 | 390 | 417 |
| 0.60 | 224 | 3397 | 623 | 655 | 226 | 3458 | 616 | 217 | 217 | 3472 | 569 | 603 | 202 | 3472 | 499 | 533 |

Table 3: Values of elastic constants at different pressures and temperatures for $\mathbf{M g B}_{2}$.

|  | 0 K |  |  |  | 300 K |  |  |  | 600 K |  |  |  | 900 K |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $V / V_{0}$ | $\begin{gathered} \hline \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ | $\begin{gathered} \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ | $\begin{gathered} \hline \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ | $\begin{gathered} \hline \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ |
| 1 | 0 | 462 | 67.0 | 80.0 | 0 | 458 | 63.0 | 76.0 | 0 | 447 | 51.7 | 64.7 | 0 | 432 | 37.4 | 50.4 |
| 0.95 | 8.58 | 562 | 81.5 | 97.3 | 8.38 | 559 | 76.9 | 92.8 | 7.75 | 548 | 63.4 | 79.3 | 6.95 | 532 | 46.0 | 62.0 |
| 0.90 | 19.7 | 688 | 99.8 | 119 | 19.2 | 687 | 94.6 | 114 | 17.8 | 676 | 78.2 | 97.8 | 16.0 | 658 | 57.0 | 76.8 |
| 0.85 | 34.0 | 846 | 123 | 147 | 33.4 | 848 | 117 | 141 | 31.0 | 837 | 96.8 | 121 | 27.9 | 817 | 70.8 | 95.3 |
| 0.80 | 52.7 | 1044 | 151 | 181 | 51.9 | 1050 | 144 | 174 | 48.4 | 1038 | 120 | 150 | 43.7 | 1017 | 88.0 | 119 |
| 0.75 | 77.3 | 1291 | 187 | 224 | 76.4 | 1300 | 179 | 216 | 71.5 | 1290 | 149 | 187 | 64.8 | 1266 | 110 | 148 |
| 0.70 | 110 | 1621 | 235 | 281 | 109 | 1637 | 225 | 272 | 103 | 1627 | 188 | 235 | 93.4 | 1601 | 139 | 187 |
| 0.65 | 154 | 2045 | 297 | 354 | 154 | 2069 | 285 | 343 | 145 | 2061 | 238 | 298 | 133 | 2031 | 176 | 237 |
| 0.60 | 215 | 2605 | 378 | 451 | 215 | 2640 | 363 | 438 | 204 | 2634 | 305 | 381 | 187 | 2601 | 225 | 303 |

Table 4: Values of elastic constants at different pressures and temperatures for $\mathbf{T i B}_{2}$.

| $V / V_{0}$ | 0 K |  |  |  | 300 K |  |  |  | 600 K |  |  |  | 900 K |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\begin{gathered} \hline \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ | $\begin{gathered} \hline \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ | $\begin{gathered} \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ | $\begin{gathered} \mathbf{P} \\ (\mathbf{G P a}) \end{gathered}$ | $C_{11}$ | $C_{12}$ | $C_{44}$ |
| 1 | 0 | 660 | 48.0 | 260 | 0 | 656 | 44.0 | 256 | 0 | 645 | 32.7 | 245 | 0 | 630 | 18.4 | 230 |
| 0.95 | 11.5 | 733 | 53.3 | 289 | 11.3 | 731 | 49.0 | 285 | 10.7 | 720 | 36.5 | 274 | 9.96 | 706 | 20.6 | 258 |
| 0.90 | 24.9 | 822 | 59.8 | 324 | 24.5 | 822 | 55.1 | 321 | 23.2 | 813 | 41.2 | 309 | 21.6 | 799 | 23.3 | 292 |
| 0.85 | 40.1 | 928 | 67.5 | 366 | 39.6 | 932 | 62.5 | 364 | 37.7 | 324 | 46.8 | 351 | 35.2 | 912 | 26.6 | 333 |
| 0.80 | 57.4 | 1055 | 76.8 | 416 | 56.8 | 1063 | 71.3 | 415 | 54.2 | 1058 | 53.7 | 402 | 50.8 | 1047 | 30.6 | 382 |
| 0.75 | 76.2 | 1208 | 87.8 | 476 | 75.8 | 1220 | 81.8 | 476 | 72.7 | 1219 | 61.8 | 463 | 68.3 | 1210 | 35.3 | 442 |
| 0.70 | 96.0 | 1403 | 102 | 553 | 96.0 | 1423 | 95.4 | 555 | 92.6 | 1427 | 72.4 | 542 | 87.5 | 1421 | 41.5 | 519 |
| 0.65 | 115 | 1643 | 119 | 647 | 116 | 1672 | 112 | 653 | 112 | 1684 | 85.4 | 640 | 107 | 1684 | 49.2 | 615 |
| 0.60 | 128 | 1946 | 142 | 767 | 131 | 1990 | 133 | 776 | 129 | 2011 | 102 | 764 | 125 | 2018 | 58.9 | 737 |

## CONCLUSION

It is found from the results that the elastic constants increase substantially with the increase in pressure along each isotherm at $0 \mathrm{~K}, 300 \mathrm{~K}, 600 \mathrm{~K}$ and 900 K . At $\mathrm{P}=0$, the elastic constants decrease with the increase in temperature. At higher compressions the change in elastic constants becomes small.

## ACKNOWLEDGEMENTS

The author is thankful to Professor Rajpal Dahiya, Centre for Energy Studies, Indian Institute of Technology, New Delhi and Dr. S.K. Rathi, Additional Director, B.S.A. College of Engineering and Technology, Mathura (U.P.) for inspiration and guidance.

## REFERENCES

1. Dhoble, A. and Verma, M.P. (1986), Thermodynamic Analysis of Anderson- Gruneisen parameter, Phys. Stat. Sol. (b) 133, 491.
2. Kumar, M. and Dass, N. (1984), High Pressure equation of State for Solids, Phys. Stat. Sol. (b) 127, 101.
3. Thomas, L.M. and Shankar, J. (1995), Equation of state and pressure dependence of Bulk Modulus of NaCl crystals, Phys. Stat. Sol. (b) 189, 363.
4. Anderson, O.L., Isaak, D.G. and Oda, H. (1992), High Temperature elastic constant data on minerals relevant to geophysics, Rev. Geophysics, 30, 57.
5. Holzapfel, W.B. (1996), Equation of State for Strong Compression, Rep. Prog. Phys. 59, 29.
6. Stacy, F.D. (2001), Finite Strain Thermodynamics and Earth's Core, Phys. Earth Planet Inter. 128, 179.
7. Birch, F. (1986), Equation of State and Thermodynamic parameters of NaCl to 300 K bar in high temperature domain, J. Geophys. Res. 42, 4949.
8. Gaurov, S., Sharma, B.S., Sharma, S.B. and Bindu, S. (2002), Elastic Properties of intermetallic compounds under high pressure \& high temperature, Physica, B 322, 328.
9. Bina, C.R., Helffrich, G..R. (1992), Calculation of Elastic Properties from Thermodynamic Equation of State Principles, Annu. Rev. Earth Planet Sci. 20, 527.
10. Upadhyay, A.K., Sharma, B.S. (2009), Pressure Dependence of Elastic Constants of Intermetallic Compunds, Acta Ciencia Indica, 25, 259.
11. Liu, Z.L., Chen, X.R. and Wang, Y.L. (2006), Electronic Structure and thermodynamic properties of LiBc under high pressure, Physica, B. 381, 139.
