

The Effect of Monte Carlo, Molecular Dynamic and Langevin Dynamic Simulation and Computational Calculations on Insulin-like Growth Factor-1(IGF-1)

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Abstract

Insulin-like growth factor (IGF-1) is an anti-apoptosis factor in multiple cell types associated with various cancers. Computational methods allow investigating the systems between 50–100 atoms in the frame of quantum mechanics and up to 50,000 atoms with molecular dynamics. Since there are specific interactions between the residues, the solvent could play an important role in the stability of the native structure. Therefore it is useful to carry out such simulations at atomistic detail. MC, MD and LD simulations of the IGF-1 were performed with the HyperChem7.0 program. The geometries, and the interaction energies, bonds, angles, stretch-bends, electrostatic and the VDW Interactions were carried out in solution and gas phase. We have computed the transition temperature for the IGF-1 molecule. Studying the changes occurred in the potential energy of the three force fields showed that Amber force field is better than MM+ and OPLS force field and also MD simulation, at least in this model, is more effective than MC and LD methods. After equilibration, the MD simulation was very stable, and the difference between the relation coefficients $R^2=0.8173$ in gas and $R^2=0.7558$ in water was compared. The Pearson correlation suggests that there is an inverse relationship ($R=-0.25$) between *in vitro* temperature and stability of the structure.

Keywords: IGF-1, AMBER, MM+, OPLS, MC, MD, LD

Introduction

Insulin-like growth factor (IGF-1), also known as somatomedin C, mediates the growth promoting activity of growth hormone. IGF-1 is autocrine regulator of cell proliferation, paracrine growth and survival factor for mammalian embryo development (Emmitte et al., 2009). Recent NMR studies have revealed that IGF-1 has three α -helical regions surrounding a hydrophobic core (Laajok et al., 2000).

Over-expression or auto-activation of the insulin-like growth factor-1 receptor (IGF-1R) tyrosine kinase has been associated with various cancers. Insulin-like growth factor (IGF-1) is an anti-apoptosis factor of multiple cell types, and the anti-apoptotic effects are mediated through mitochondrial and cytochrome-c pathway (Li et al., 2003).

Development of faster computers that are within

the reach of the widest scientific community as well as efficient computational methods allows investigating systems between 50–100 atoms in the frame of quantum mechanics and up to 50,000 atoms with molecular dynamics. Since the models become increasingly realistic, direct comparison with experimental data becomes possible (Na'ray-Szabo and Berenteb, 2003). In addition to hit identification, docking techniques are increasingly used to support lead optimization efforts (Kitchen et al., 2004).

Recently, constant temperature molecular simulations of peptide folding have been reported using implicit solvent models, and explicit solvent models (Sung and Wu, 1997; Daura et al., 1998)

Recently, however, several computer simulations have demonstrated a strong coupling between hydrophobicity, solute-solvent dispersion attractions, and electrostatics. For example, simulations of explicit water between plate like solutes revealed that hydrophobic attraction and dewetting phenomena are strongly sensitive to the nature of solute-solvent dispersion interactions (Dzubiell, 2006). The competing effects of the

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solvent, such as the Van Der Waals (VDW) attraction and hydrogen bonding between the protein and solvent, reduce the strength of the interactions and consequently reduce the energy barrier related to the multiple minima problem (Ozkan, 2004).

In a solution, the intramolecular VDW interactions of a protein molecule are balanced by the intermolecular VDW interactions with solvent molecules. The possible difference between the protein intramolecular VDW attraction and that with water may be included in the hydrophobic interaction energy (Sung, 1999). Kurochkina and Lee have shown that the pair-wise sum of the buried surface area is linearly related to the true buried area (Sung, 1999).

Since the specific interactions between the residues and solvent play an important role in the stability of the native structure, it is useful to carry out such simulations at atomistic detail. This comes with the problem of timescale of folding/unfolding that is several orders of magnitude larger than those currently attainable by MD simulations (Ozkan, 2004).

Water plays a crucial role for the stability, dynamics, and function of proteins. For this reason, Molecular Dynamics (MD), Monte Carlo (MC) and Langevin Dynamic (LD) simulations must account for the effects that this solvent has, both on protein structure and on protein dynamics (Hamaneh and Buck, 2007).

The research group used structural bioinformatic methods to study the structure, function and dynamics of biomolecules. Our goal was to better understand structure formation of biomolecules and their mechanism. The function of proteins or nucleic acids depends on its three-dimensional structure and conformational dynamics. We employed the molecular dynamics simulation method as the main tool to study conformational dynamics of biomolecules. This approach allowed following dynamic molecular processes at high resolution in space and time. It was also possible to study the thermodynamics and energetics of structure formation and association. Our research was not restricted to a single biomolecular system.

The aim of the present work was to describe and characterize the molecular structure of IGF-1 crystalline-structure. In this work, the structure of a coordination compound modeling the IGF-1 was discussed computationally. Thus, it was worthwhile to collect information on its structure by means of computational chemistry as well.

Materials and Methods

Monte Carlo simulations are based on pair wise additive potentials of the form (Tafazzoli and Khanlarkhani, 2007). In concepts and algorithms of classical MD simulations the atoms of a biopolymer move according to the Newtonian equations of motion (Phillips et al., 2005)

For an all-atom MD simulation, one assumes that every atom experiences a force specified by a model force field accounting for the interaction of that atom with the rest of the system.

$$E_{\text{total}}=E_{\text{bond}}+E_{\text{angle}}+E_{\text{dihedral}}+E_{\text{vdw}}+E_{\text{coulomb}} \quad (1)$$

The Van Der Waals potential energy for the general treatment of non-bonded interactions is often modeled by a Lennard–Jones 12–6 function. We can consider an effective Hamiltonian operator constructed for molecule in a given geometry of it and the solvent:

$$H_{\text{eff}}=H_0+V_{\text{elec}}+V_{\text{ind}}+V_{\text{non-elec}}, \quad (2)$$

Where H_0 is the Hamiltonian in gas phase (the unperturbed Hamiltonian), V_{elec} is the perturbation from the permanent charge distribution of water, represented as a set of point-charges, V_{ind} is the perturbation from the induced dipoles in the solvent and $V_{\text{non-elec}}$ is a non-electrostatic perturbation, which models the effect of the anti-symmetry between the solute and solvent (Hermida-Ramón et al., 2009)

Using Langevin dynamics, solvent effects can be modeled and the dynamic behavior of a molecular system in a liquid environment can be studied. These simulations can be much faster than molecular dynamics. These simulations can be used to study the same kinds of problems as molecular dynamics: time dependent properties of solvated systems at non-zero temperatures. Because of the implicit treatment of the solvent, this method is particularly well-suited for studying large molecules in solution.

Langevin dynamics simulates the effect of molecular collisions and the resulting dissipation of energy that occurs in real solvents, without explicitly including solvent molecules. This is accomplished by adding a random force and a frictional force to each atom at each time step (Berendsen, 1990; Karplus and Petsko 1990).

The random force is calculated as a random number, taken from a Gaussian distribution with a mean value of zero and with no correlation with the atom's velocity.

Molecular mechanics (MM) force fields rely on the combination of Coulomb and Lennard–Jones interactions to describe all non-bonded interactions (Ponder and Case, 2003). Even though the functional form of the potential energy is quite simple, it depends on a large number of empirical parameters, which must be obtained from *ab initio* calculations of the electronic structure in small molecules and/or experimental data.

Because each new term in the MM potential function requires additional empirical parameters, it is quite appealing to keep the functional form of the potential function as simple as possible. While most widely-used current force fields such as AMBER, OPLS do not employ explicit hydrogen bonding terms, this was not always the case (MacKerrell et al., 1998; Hagler and Lifson, 1974; Cornell et al., 1995; Jorgensen et al., 1996; Weiner et al., 1984).

The crystal structures of proteins were from the Brookhaven Protein Data Bank. The structure of protein IGF-1 was selected from the Protein Data

Bank (PDB code 1B9G). These studies provided insights into the steric, electrostatic, hydrophobic, and hydrogen bonding properties and other structural features influencing the IGF-1.

In vacuum, the system was simulated using Monte Carlo, Molecular dynamic and Langevin dynamics with 100 ps step and without any constraints. Temperature was kept constant at 300 K. In water, simulations, the system was placed in a box (3 x 3 x 3 nm) containing one molecule of solute and 884 TIP3P water molecules (figure 1). The system was simulated using Newtonian dynamics with 100 ps step and no constraints applied to the solute.

The complex was solvated by added water molecules. The systems were first energy minimized steps with the conjugate gradient algorithm. Then, the position-restrained MC, MD and LD simulation were run 100ps. Afterwards, 1 ps simulations were carried out at a time step of 100ps (figure 1).

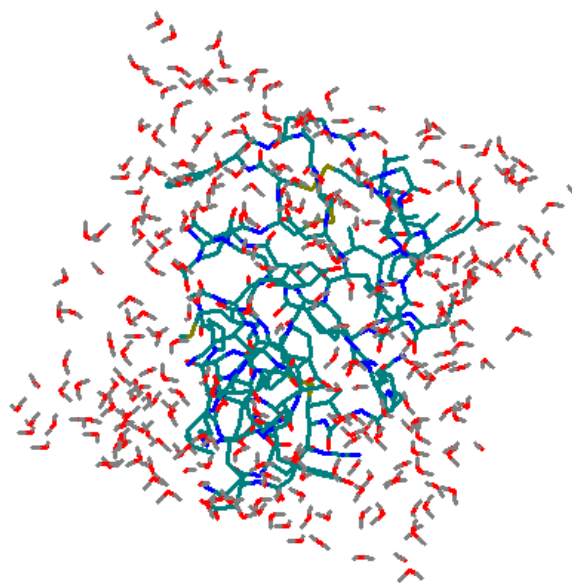


Figure 1. Schematic representation of structural model of IGF-1R in water (884 TIP3P water molecules).

Several simulations were carried out, as listed in Table 1. MC, MD and LD simulations of the IGF-1 were performed with the HyperChem7.0 program (HyperChem 7.0, 2001). The geometries, and the interaction energies, bonds, angles, stretch-bends, electrostatic and the VDW interactions were carried out in solution and gas phase (table 1, figure 2). In the simulation of the small water system, with temperature of 295, 297, 299, 301, 303 and 305 K, the instantaneous kinetic temperature is given by Eq (3), (Español and Warren, 1995):

$$T_k(t) = \sum_{i=1}^N m_i v_i^2(t) / k_B N_f \quad (3)$$

Where k_B is the Boltzmann constant, N_f is the degrees of freedom ($N_f = 3N - 3$ for a system of N particles with fixed total momentum), m_i is the atom weight for atom i , v_i is the velocity of atom i . The effect of confinement on the thermodynamic properties of several statement proteins was investigated by performing simulations over a large range of temperatures. We have computed the transition temperature for the IGF-1 molecule.

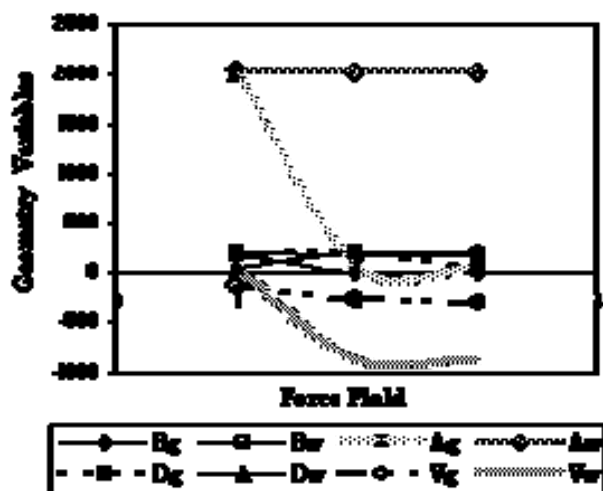


Figure 2. Geometry of optimized variables of Bond length (B), Bond Angle (A) and Dihedral Angle (D) in gas and water media at 300K.

Table 1. Calculation of various variables in 300 K Temperature for IGF-1 at MM+, AMBER, OPLS.

environment	GEOMETRY						MC	MD			LD		
	Force Field	Bond	Angle	Dihedral	Energy	Gradient	Potential	Potential	Kinetic	Total Energy	Potential	Kinetic	Total Energy
gas	MM+	174.408	2003.05	187.056	2186.66284	0.099337	547.949	2456.24	386.747	2842.99	354.473	388.198	742.671
	AMBER	11.0393	67.3022	200.754	23.23151	0.099951	2798.14	348.117	393.388	741.505	2481.28	394.672	2875.95
	OPLS	2.53763	47.7753	48.4863	-202.51639	0.098917	349.788	138.21	390.617	528.827	131.714	391.538	523.253
water	MM+	202.777	2037.84	43.0804	-130.11968	0.095515	2217.07	6677.62	9695.98	16373.6	672.07	1562	2300.2
	AMBER	175.662	2019.97	212.994	270.345123	0.084089	160.52	733.188	1608.57	2341.76	746.05	1596	2342.05
	OPLS	175.662	2019.97	212.994	270.345123	0.084089	1367.75	686.563	1826.73	2513.3	641.945	1871.01	2512.95

Results

The following text describes methods for generating and evaluating representative molecular conformations, particularly for peptides and small proteins, based on Molecular Mechanic energy functions. On the other hand, Molecular Mechanics describe molecules as atoms linked with springs (harmonic bond stretches and bond angle wagging), each atom having finite volume and relatively sharp boundaries ("6-12" hard spheres potentials), with sinusoidal torsion energies. The force field for a typical protein can be given as a sum of the various components including bond stretching and bending, torsion potentials, and non-bonded interactions.

Molecular dynamic simulations have been widely used to obtain the 'real' bioactive conformation when the crystal structure of protein–ligand complex is unavailable. So, in order to obtain the 'real' stabilized bioactive compound was used for molecular dynamic simulations. The system was well-equilibrated and 500 ps in the range of the MD equilibration were selected for

further processing analysis. We have shown relation coefficients in figure 3.

In this paper, we have used Monte Carlo methods to study IGF-1 in the bulk and in a confined environment. Results are presented in table 1 in respect of the effects on the specific media of the structure.

Simulations of molecular dynamics were carried out on the two systems, gas and solvent IGF-1 molecules. All simulations were carried out at the constant temperature. All simulations were performed at 300 K. Each solvent system was immersed in a periodic water box, and the structures of water molecules were maintained. A 100 ps time step was used in all simulations.

Furthermore, we used MD and LD methods to study protein in the bulk and in confined media. The structures obtained throughout this calculation were optimized using MM+, AMBER, and OPLS force field parameters. Also, all these approaches included discrete particles moving in a defined energy landscape according to Langevin Dynamics (LD).

K. For example, by the iterative calculation with time step of "100 ps", using "Chem3D" software, the experimental X-ray geometrical data reported for IGF-1 in crystalline structure, were used as input geometrical data (Hartung et al., 1992). At the next step, the appropriate ones, carefully selected

from the structures obtained throughout this calculation, were optimized using MM+, AMBER and OPLS force field parameters and included into the same software. In this paper, a comprehensive conformational research on free molecules was carried out.

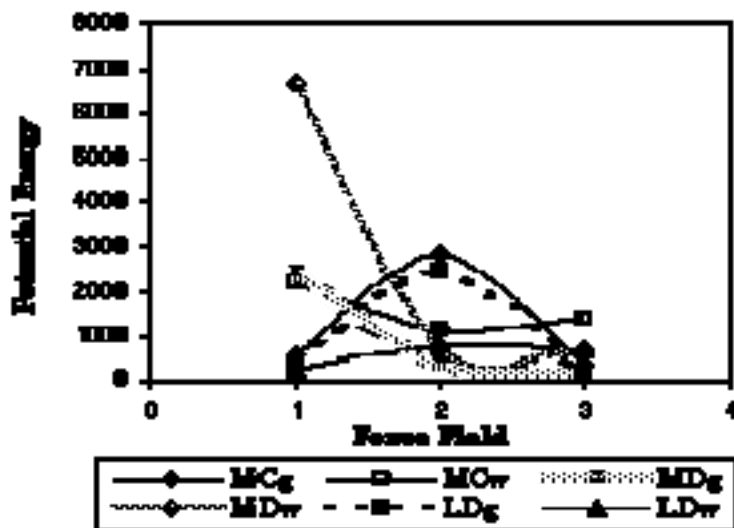


Figure 3. The potential energy (kcal/mol) via time (ps) during Molecular Dynamic (MD) simulation at 300K in gas ($R2=0.8173$) and water ($R2=0.7558$) environments to a stabilized structure of IGF-1.

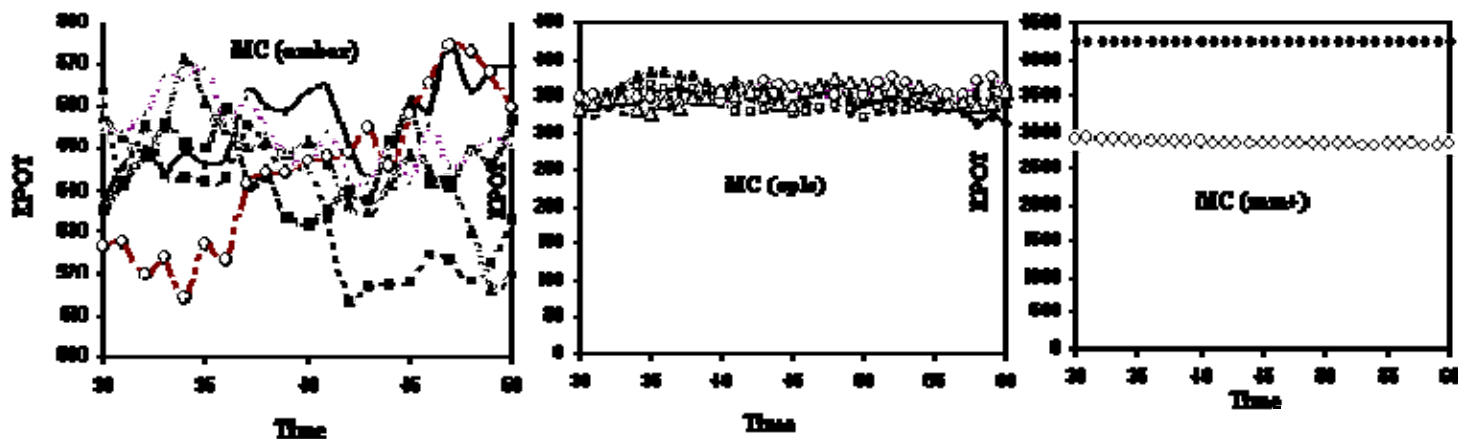


Figure 4. The potential energy(kcal/mol) via time (ps) during Monte Carlo (MC) simulation at 295,297,299,301,303 and 305 K using (a) AMBER (b) OPLS (c) MM+ force fields corresponding to a stabilized structure of IGF-1.

Potential energies for the three force fields of MM+, AMBER and OPLS at Monte Carlo simulation were compared in figure 4. The average energies are in good agreement with the simulation accuracy. The sampling results of step-size of MD and LD methods are presented in figures 5 and 6, respectively.

Discussion

In conclusion, in this work we used molecular dynamic models to explore the stability of IGF-1 by comparing theoretical methods of simulation. A highly selective effect of temperature and environment was discovered in chemical structure and it investigated the standard constant

temperature at MC, MD and LD simulations. Molecular Dynamic simulation method has been employed as the main tool for studying conformational dynamics of biomolecules.

One of the force fields designed for treating macromolecules can be simplified by not considering explicitly – the so-called united atom approach is AMBER. Seemingly, solvent effects influence the calculated potential energy surface by lowering potential energy barriers on angle. This means that the parameterizations that have been developed for small molecules with considerable effort can be carried over into macromolecular calculations with little or no change. Also, the MM+ and OPLS force fields parameters were applied for IGF-1 model in gas and for water environment.

Also, the possible difference between the IGF-1 intramolecular VDW attraction and that with water were included in the hydrophobic interaction energy. The short-range repulsion represents the exclusive volume of each atom and needs to be calculated explicitly.

In a solution, the intramolecular VDW interactions of a protein molecule are balanced by

the intermolecular VDW interactions with solvent molecules. Thus, when solvent molecules are not explicitly included, the intramolecular VDW interactions must be adjusted accordingly. The longer-range attractive VDW interactions provide a nearly uniform background potential (Chandler et al., 1983), and therefore can serve as the reference for the VDW energy calculation (McCammon et al. 1980).

The measurement of the potential of solvation under similar conditions of temperature in solution along with investigation of energetic and structural aspects of solution were used to gain insight into the molecular level interaction with IGF-1. Solute-solvent pair interaction of potential energies show that the greater stability of solvent observed over all states investigated in this study is related to the MD/AMBER approach. The obtained results have demonstrated that the free molecule has a very flexible macro-cyclic structure. On the basis of the theoretical results obtained for the determined most stable, the dependencies of the geometrical and force constants parameters of the free molecule to its conformational structure were discussed.

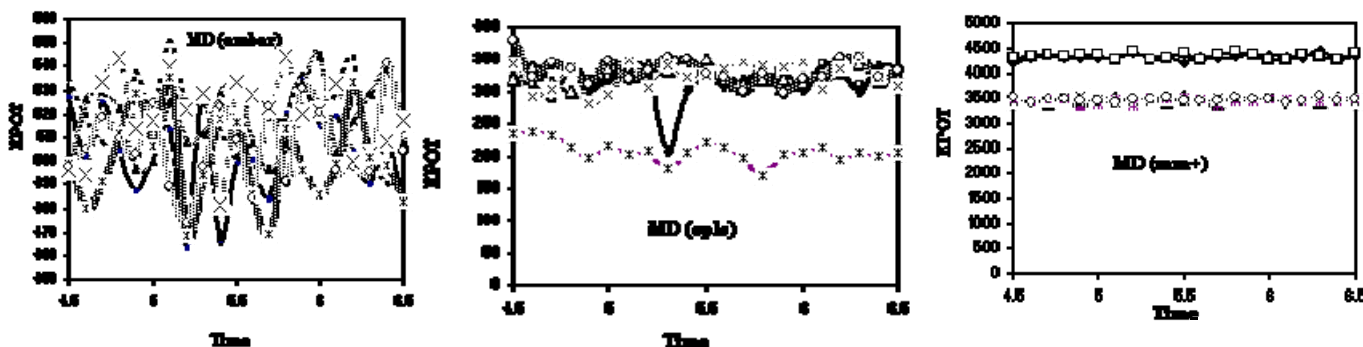


Figure 5. The potential energy (kcal/mol) via time (ps) during Molecular Dynamic (MD) simulation at 295,297,299,301,303 and 305 K using a) AMBER b) OPLS and c) MM+ force fields corresponding to a stabilized structure of IGF-1.

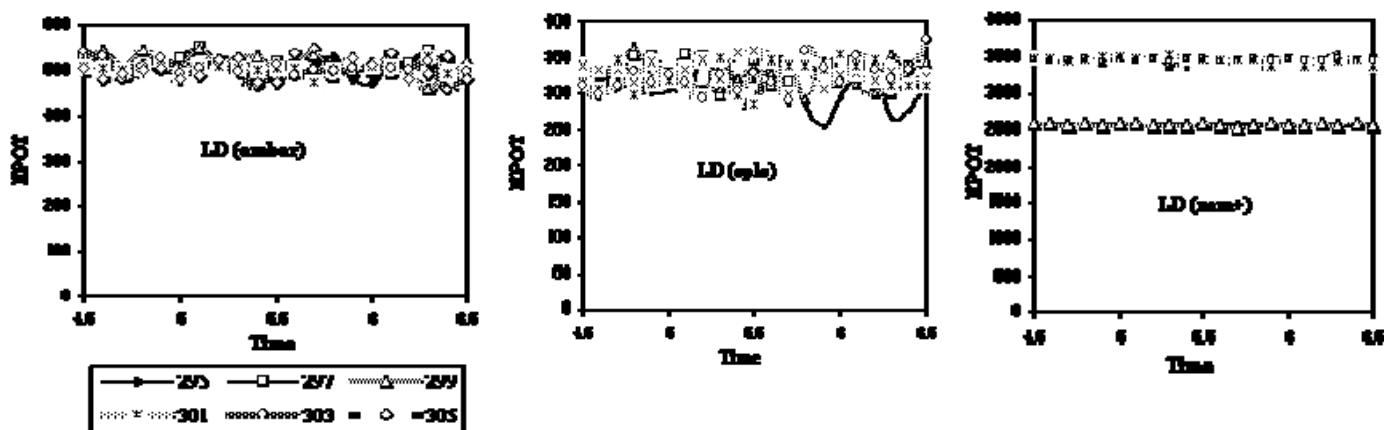


Figure 6. The potential energy (kcal/mol) via time (ps) during Langevin Dynamic (LD) simulation at 295,297,299,301,303 and 305 K using a) AMBER b) OPLS and c) MM+ force fields corresponding to a stabilized structure of IGF-1.

The potential energy, represented through the MD “force field” is the most crucial part of the simulation since it must faithfully represent the interaction between atoms, which is to be cast in the form of a simple mathematical function that can be calculated quickly. After equilibration, the MD simulation was very stable; therefore we tried to compare the difference between the relation coefficients $R^2=0.8173$ in gas and $R^2=0.7558$ in water. These results show that the force field of AMBER has a convenient relation in all simulation methods as well as in various media (figure 3).

As expected, AMBER demonstrates much smoother energy profiles than the other two simulation methods due to higher-order energy conservation in the modified Hamiltonian (figure 4a). The magnitudes of energy fluctuations in both MM+ and OPLS approaches are significantly smaller than the other (figure 4b, c). Observed data are almost identical for both choices of the MD simulation length suggesting that the MD simulation had been affected much more on acceptance rate - at least for this particular model - than MC and LD approaches. This potential does not have any terms describing angular dependencies of hydrogen bonds and is similar to the 10–12 hydrogen bonding potential originally proposed by others (McGuire et al., 1972). They found that hydrogen bonding energies were represented adequately by a sum of Lennard–Jones and electrostatic interactions plus the 10–12 hydrogen bonding term with empirical constants adjusted according to the hydrogen bond type.

Because the functional form of such a hydrogen bonding term was very close to the Lennard–Jones component of the force field, the second-generation AMBER force field omitted it altogether (Cornell et al., 1995), relying instead on the combination of Lennard–Jones and Coulomb interactions to model hydrogen bonded complexes, thus the data of this force field in three simulation methods showed that the changes of potential energy via time at various temperatures are much better than MM+ and OPLS force fields (figure 4a, 5a, 6a).

Similarly, the widely used OPLS force field does not contain an explicit hydrogen bonding term: the emphasis of OPLS parameterization is on reproducing thermodynamic properties of organic liquids such as enthalpies of vaporization, densities and free energies of hydration (figure 4b, 5b, 6b) (Jorgensen, 1996; Jorgensen and Tirado-Rives, 1988).

Because each new term in the MM+ potential function requires additional empirical parameters, it is quite appealing to keep the functional form of the potential function as simple as possible (figure 4c,

5c, 6c).

MD simulation is suitable for obtaining the elastic properties of a system in our size. For certain confining environments, individual proteins do exhibit power-law dependence, but the relationship is different for each molecule. In other cases, the increase in stability upon confinement interestingly demonstrates no monotonic behavior. ANOVA showed that there is a significant difference between the differences in the potential energy in the amber force field and MM+ and OPLS (P -value=0.035). Furthermore, the Pearson correlation suggests that there is an inverse relationship ($R=-0.25$) between in-vitro temperature and stability of the structure.

These results also revealed that the solvation of IGF-1 is the major component for the interaction potential energy and it was clearly shown that the role of the solute-solvent interactions is more pronounced in IGF-1 solvation. The major part of this difference is due to the interaction of IGF-1 with solvent molecules corresponded to various simulation methods and force fields. A difficult task in the computational study of stabilized structure is to find a proper energy function that can lead to a unique structure. Our simulations showed that the simple energy function modified to include solvent effect has a parameter range that can simulate indicated structure at constant temperature of 300K.

Several molecular dynamic simulations could be performed over a wide range of temperatures, and the data could be combined using a weighted histogram approach (Weiner et al., 1984); however, the statistical error associated with the tails of the sampled distributions is usually large and can propagate when data from simulations at different temperatures are merged.

From the simulations, it is demonstrated that the kinetic temperature of the system is properly bounded around the prescribed equilibrium temperature. The length of each simulation was 100 ps. The relative drift of molecular temperature denoted by ΔT in percent, with respect to mean temperature; T in Kelvin has been measured. Figures showed the function of the reduced temperature. Low reduced temperatures promote complex structure stability, whereas high reduced temperatures oppose it.

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